# PERTURBED NONLINEAR ELLIPTIC NEUMANN PROBLEMS INVOLVING ANISOTROPIC SOBOLEV SPACES WITH VARIABLE EXPONENTS 

A. AHMED - M.S.B. ELEMINE VALL

In this paper we study the existence of infinitely many weak solutions of the following perturbed Kirchhoff-type non-homogeneous Neumann problem

$$
\left\{\begin{array}{l}
-\sum_{i=1}^{N} M_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}}\right) \\
+M_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|u|^{p_{0}(x)}\right)|u|^{p_{0}(x)-2} u=f(x, u)+g(x, u) \text { in } \Omega \\
\sum_{i=1}^{N}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}} v_{i}=0 \quad \text { on } \quad \partial \Omega
\end{array}\right.
$$

by applying technical approach based on critical points theorem due to B. Ricceri in a reflexive anisotropic Sobolev spaces. We use some suitable assumptions on the right had side but without using log-Hölder continuous condition.

## 1. Introduction

In recent years, the anisotropic variable exponent Sobolev spaces have attracted the attention of many mathematicians, physicists and engineers. The impulse

## Received on December 25, 2022

AMS 2010 Subject Classification: 35R11, 35A15, 35J60, 35D05, 35D30, 35J62.
Keywords: Ricceri's variational principle, Kirchhoff-type problem, Non-homogeneous operators, Elliptic problems, anisotropic variable exponent Lebesgue-Sobolev spaces.
for this mainly come from their important applications in modelling real-world problems in electrorheological, magnetorheological fluids, elastic materials and image restoration, (see for example [9], [15], [16], [35], [38], [39]).

More recently, several authors (see e.g. [2], [10], [25]) have studied the anisotrop-ic quasi-linear elliptic equations with variable exponents, i.e. the quasi-linear elliptic equations involving the following $\vec{p}(\cdot)$-Laplacian

$$
\begin{equation*}
\Delta_{\vec{p}(\cdot)} u=\sum_{i=1}^{N} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}}\right) \tag{1}
\end{equation*}
$$

It's clear that this $\vec{p}(\cdot)$-Laplace operator is a generalization of the $p(\cdot)$-Laplace operator

$$
\begin{equation*}
\Delta_{p(\cdot)} u=\operatorname{div}\left(|\nabla u|^{p(x)-2} \nabla u\right) \tag{2}
\end{equation*}
$$

We refer to [1], [37], [40] for the study of the $p(\cdot)$-Laplacian equations and the corresponding variational problems.

The $p(\cdot)$-Laplacian is a meaningful generalization of the $p$-Laplacian operator

$$
\begin{equation*}
\Delta_{p} u=\operatorname{div}\left(|\nabla u|^{p-2} \nabla u\right) \tag{3}
\end{equation*}
$$

obtained in the case when $p$ is a positive constant.
On the one hand, Ricceri [33], Anello and Cordaro [8] studied the existence of solutions for the following problem

$$
\begin{cases}-\Delta_{p} u+a(x)|u|^{p-2} u=b(x) f(u)+c(x) g(u) & \text { in } \Omega  \tag{4}\\ \frac{\partial u}{\partial v}=0 & \text { on } \partial \Omega\end{cases}
$$

where $a(x)$ is a positive function such that $a(\cdot) \in L^{\infty}(\Omega)$ with $a^{-}=\operatorname{essinf}_{x \in \Omega} a(x)>0$ and $p>N$. The existence of solutions of problem (4) was proved by applying Ricceri's variational principle (see [32]).

In [21], X. Fan, C. Ji treated the problem

$$
\begin{cases}-\Delta_{p(\cdot)} u+a(x)|u|^{p(x)-2} u=f(x, u)+g(x, u) & \text { in } \Omega  \tag{5}\\ \frac{\partial u}{\partial v}=0 & \text { on } \partial \Omega\end{cases}
$$

and they proved the existence of infinitely many solutions in the variable exponent Sobolev space $W^{1, p(\cdot)}(\Omega)$.

However, there are some non-homogeneous materials that have different behaviors in different space directions, hence the need for anisotropic spaces
with variable exponent. Ahmed, Hjiaj, and Touzani have studied in [3] the Neumann $\vec{p}(\cdot)$-elliptic problem :

$$
\begin{cases}-\Delta_{\vec{p}(\cdot)} u+a(x)|u|^{p_{0}(x)-2} u=f(x, u)+g(x, u) & \text { in } \Omega  \tag{6}\\ \frac{\partial u}{\partial v}=0 & \text { on } \partial \Omega\end{cases}
$$

and they proved the existence of infinitely many weak solutions in the anisotropic variable exponent Sobolev space $W^{1, \vec{p}(\cdot)}(\Omega)$ under some hypotheses. For other related results, we refer to [5], [6], [14], [19].

On the other hand, much interest has been focused on the study of Kirchhoff type problems. More precisely, Kirchhoff studied the following model problem (see [26]) as an extension of d'Alembert's classical wave equation by considering changes in string length during vibrations

$$
\begin{equation*}
\rho \frac{\partial^{2} u}{\partial t^{2}}-\left(\frac{P_{0}}{h}+\frac{E}{2 L} \int_{0}^{L}\left|\frac{\partial u}{\partial x}\right|^{2} d x\right) \frac{\partial^{2} u}{\partial x^{2}} \tag{7}
\end{equation*}
$$

where $L$ is the length of the chord, $h$ is the area of the cross section, $E$ is the Young's modulus of the material, is the density and $P_{0}$ is the initial tension. A distinguishing feature of the Kirchhoff equation (7) is that the equation contains a non-local coefficient $\frac{P_{0}}{h}+\frac{E}{2 L} \int_{0}^{L}\left|\frac{\partial u}{\partial x}\right|^{2} d x$ which depends on the average $\frac{E}{2 L} \int_{0}^{L}\left|\frac{\partial u}{\partial x}\right|^{2}$ of the kinetic energy $\frac{1}{2}\left|\frac{\partial u}{\partial x}\right|^{2}$ on $[0, L]$, and hence the equation is no longer a point-wise identity. See also [11], [17], [23], [34], [36] for related topics.

The purpose of our paper is to investigate a class of Kirchhoff type problems involving operators in divergence form as follows:

$$
\begin{cases}-\sum_{i=1}^{N} M_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}}\right) &  \tag{8}\\ +M_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|u|^{p_{0}(x)}\right)|u|^{p_{0}(x)-2} u=f(x, u)+g(x, u) & \text { in } \Omega \\ \sum_{i=1}^{N}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}} v_{i}=0 & \text { on } \quad \partial \Omega\end{cases}
$$

where $M_{i}, f$ and $g$ define and satisfies some conditions detailed in Section 3.
In the Dirichlet case, A. Ourraoui in [29] have studied the problem (8) by assuming an Ambrosetti-Rabinowitz type condition and using techniques related to a Mountain pass theorem in the case where $g \equiv 0$. M. Avci, R. A. Mashiyev and B. Cekic in [12] have studied the same problem in the case where $g \equiv 0$,
$M_{0}=M_{2}=\cdots=M_{N}=M$ and the assumption $2 \leq p_{i}(x) \leq N$. Note that the hypotheses we adopt are totally different from the ones assumed in the papers just quoted.

It is no a surprise that the presence of Neumann conditions in an anisotropic non-homogeneous perturbed-Kirchhoff type problem make difficulties in the application of the Theorem 2.3 which is our main tool. To overpass these difficulties, we combine the classical techniques with the recent techniques that appeared when treating anisotropic problems with variable exponents.

This paper is organized as follows. In Section 2, we recall some basic facts about anisotropic variable exponent Sobolev spaces as well as Ricceri's variational principle. In Section 3, we state and prove our main results (Theorems 3.5 and 3.6 ), providing also some remarks and examples.

## 2. Preliminary results

Let $\Omega$ be a bounded domain in $\mathbb{R}^{N}$, we define:
$\mathcal{C}_{+}(\bar{\Omega})=\left\{\right.$ measurable function, $p(\cdot): \bar{\Omega} \longrightarrow \mathbb{R} \quad$ such that $\left.\quad 1<p^{-} \leq p^{+}<\infty\right\}$ where

$$
p^{-}=\operatorname{essinf}\{p(x): x \in \bar{\Omega}\} \quad \text { and } \quad p^{+}=\operatorname{ess} \sup \{p(x): x \in \bar{\Omega}\}
$$

We define the Lebesgue space with variable exponent $L^{p(\cdot)}(\Omega)$ as the set of all measurable functions $u: \Omega \longmapsto \mathbb{R}$ for which the convex modular

$$
\rho_{p(\cdot)}(u):=\int_{\Omega}|u|^{p(x)} d x
$$

is finite, then

$$
\|u\|_{L^{p(\cdot)}(\Omega)}=\|u\|_{p(\cdot)}=\inf \left\{\lambda>0: \rho_{p(\cdot)}(u / \lambda) \leq 1\right\}
$$

defines a norm in $L^{p(\cdot)}(\Omega)$ called the Luxemburg norm. The space $\left(L^{p(\cdot)}(\Omega), \|\right.$. $\left.\|_{p(\cdot)}\right)$ is a separable Banach space. Moreover, the space $L^{p(\cdot)}(\Omega)$ is uniformly convex, hence reflexive, and its dual space is isomorphic to $L^{p^{\prime}(\cdot)}(\Omega)$, where $\frac{1}{p(\cdot)}+\frac{1}{p^{\prime}(\cdot)}=1$. Finally, we have the following Hölder type inequality

$$
\begin{equation*}
\left|\int_{\Omega} u(x) v(x) d x\right| \leq\left(\frac{1}{p^{-}}+\frac{1}{\left(p^{-}\right)^{\prime}}\right)\|u\|_{p(\cdot)}\|v\|_{p^{\prime}(\cdot)} \tag{9}
\end{equation*}
$$

for all $u \in L^{p(\cdot)}(\Omega)$ and $v \in L^{p^{\prime}(\cdot)}(\Omega)$.
An important role in manipulating the generalized Lebesgue spaces is played by the modular $\rho_{p(\cdot)}$ of the space $L^{p(\cdot)}(\Omega)$. We give the following result.

Proposition 2.1. (See [18], [24]). If $u \in L^{p(\cdot)}(\Omega)$, then the following properties hold true:
(i) $\|u\|_{p(\cdot)}<1 \quad($ respectively $,=1,>1) \quad \Leftrightarrow \quad \rho_{p(\cdot)}(u)<1 \quad$ (respectively,$=$ $1,>1)$,
(ii) $\|u\|_{p(\cdot)}>1 \Rightarrow\|u\|_{p(\cdot)}^{p^{-}}<\rho_{p(\cdot)}(u)<\|u\|_{p(\cdot)}^{p^{+}}$,
(iii) $\|u\|_{p(\cdot)}<1 \Rightarrow\|u\|_{p(\cdot)}^{p^{+}}<\rho_{p(\cdot)}(u)<\|u\|_{p(\cdot)}^{p^{-}}$.

Now, we define the Sobolev space with variable exponent by

$$
W^{1, p(\cdot)}(\Omega)=\left\{u \in L^{p(\cdot)}(\Omega) \quad \text { and } \quad|\nabla u| \in L^{p(\cdot)}(\Omega)\right\}
$$

equipped with the following norm

$$
\|u\|_{W^{1, p(\cdot)}(\Omega)}=\|u\|_{1, p(\cdot)}=\|u\|_{p(\cdot)}+\|\nabla u\|_{p(\cdot)}
$$

The space $\left(W^{1, p(\cdot)}(\Omega),\|\cdot\|_{1, p(\cdot)}\right)$ is a separable and reflexive Banach space. We refer to [18] for the elementary properties of these spaces.
Now, we present the anisotropic Sobolev space with variable exponent which is used for the study of our main problem.

Let $p_{0}(\cdot), p_{1}(\cdot), \ldots, p_{N}(\cdot)$ be $N+1$ variable exponents in $\mathcal{C}_{+}(\bar{\Omega})$. We denote

$$
\vec{p}(\cdot)=\left\{p_{0}(\cdot), p_{1}(\cdot), \ldots, p_{N}(\cdot)\right\}, D^{0} u=u \text { and } D^{i} u=\frac{\partial u}{\partial x_{i}} \quad \text { for } i=1, \ldots, N
$$

and for all $x \in \bar{\Omega}$ we put

$$
\begin{aligned}
p_{M}(\cdot) & =\max \left\{p_{0}(\cdot), p_{1}(\cdot), \ldots, p_{N}(\cdot)\right\} \\
p_{m}(\cdot) & =\min \left\{p_{0}(\cdot), p_{1}(\cdot), \ldots, p_{N}(\cdot)\right\}
\end{aligned}
$$

We define

$$
\begin{equation*}
\underline{p}=\min \left\{p_{0}^{-}, p_{1}^{-}, p_{2}^{-}, \ldots, p_{N}^{-}\right\} \quad \text { then } \quad \underline{p}>1 \tag{10}
\end{equation*}
$$

and

$$
\begin{equation*}
\bar{p}=\max \left\{p_{0}^{+}, p_{1}^{+}, p_{2}^{+}, \ldots, p_{N}^{+}\right\} \tag{11}
\end{equation*}
$$

The anisotropic variable exponent Sobolev space $W^{1, \vec{p}(\cdot)}(\Omega)$ is defined as follows

$$
W^{1, \vec{p}(\cdot)}(\Omega)=\left\{u \in L^{p_{0}(\cdot)}(\Omega): D^{i} u \in L^{p_{i}(\cdot)}(\Omega) \quad \text { for all } i=1,2, \ldots, N\right\}
$$

endowed with the norm

$$
\begin{equation*}
\|u\|_{W^{1, \vec{p} \cdot()}(\Omega)}=\|u\|_{1, \vec{p}(\cdot)}=\|u\|_{L^{p_{0}(\cdot)}(\Omega)}+\sum_{i=1}^{N}\left\|D^{i} u\right\|_{L^{p_{i}(\cdot)}(\Omega)}, \tag{12}
\end{equation*}
$$

we refer to [13], [30], [31] for the constant exponent case.
We emphasize that the space $\left(W^{1, \vec{p} \cdot()}(\Omega),\|\cdot\|_{1, \vec{p} \cdot(\cdot)}\right)$ is a reflexive Banach space (see [20]).
A more complete theory of anisotropic variable exponent Sobolev spaces may be obtained in [4], [7], [22], [27], [28].

Throughout this paper we assume that

$$
\begin{equation*}
\underline{p}>N . \tag{13}
\end{equation*}
$$

Remark 2.2. Since $W^{1, \vec{p} \cdot \cdot)}(\Omega)$ is continuously embedded in $W^{1, \underline{p}}(\Omega)$ and $W^{1, \underline{p}}(\Omega)$ is compactly embedded in $C^{0}(\bar{\Omega})$ (the space of continuous functions), thus $W^{1, \vec{p} \cdot()}(\Omega)$ is compactly embedded in $C^{0}(\bar{\Omega})$.

Set

$$
\begin{equation*}
C_{0}=\sup _{u \in W^{1, p(r)}(\Omega) \backslash\{0\}} \frac{\|u\|_{L^{\infty}(\Omega)}}{\|u\|_{1, \vec{p}(\cdot)}} . \tag{14}
\end{equation*}
$$

Then $C_{0}$ is a positive constant.
Let us recall the following theorem obtained by B. Ricceri in [21], which will be applied to establish the existence of weak solutions for our main problem.
Theorem 2.3. (See [21], Theorem 2.2). Let E be a reflexive real Banach space, and let $\Phi, \Psi: E \longrightarrow \mathbb{R}$ be two sequentially weakly lower semi-continuous and Gâteaux differentiable functionals. Assume also that $\Psi$ is (strongly) continuous and satisfies $\lim _{\|u\|_{E} \rightarrow \infty} \Psi(u)=+\infty$. For each $\rho>\inf _{E} \Psi$, put

$$
\begin{equation*}
\varphi(\rho)=\inf _{\left.u \in \Psi^{-1}(]-\infty, \rho \mid\right)} \frac{\Phi(u)-\inf _{v \in \frac{\left(\Psi^{-1}(]-\infty, \rho \mid\right)_{w}}{}} \Phi(v)}{\rho-\Psi(u)}, \tag{15}
\end{equation*}
$$

where $\overline{\left(\Psi^{-1}(]-\infty, \rho[)\right)_{w}}$ is the closure of $\Psi^{-1}(]-\infty, \rho[)$ for the weak topology. Then, the following conclusions hold
(a) If there exist $\rho_{0}>\inf _{E} \Psi$ and $u_{0} \in E$ such that

$$
\begin{equation*}
\Psi\left(u_{0}\right)<\rho_{0}, \tag{16}
\end{equation*}
$$

and

$$
\begin{equation*}
\Phi\left(u_{0}\right)-\inf _{\left.v \in\left(\Psi-1(]-\infty, \rho_{0} \mid\right)\right)_{w}} \Phi(v)<\rho_{0}-\Psi\left(u_{0}\right), \tag{17}
\end{equation*}
$$

then the restriction of $\Psi+\Phi$ to $\Psi^{-1}(]-\infty, \rho_{0}[)$ has a global minimum.
(b) If there exists a sequence $\left(r_{n}\right)_{n} \subset\left(\inf _{E} \Psi,+\infty\right)$ with $r_{n} \rightarrow \infty$ and a sequence $\left(u_{n}\right)_{n} \subset E$ such that for each $n$

$$
\begin{equation*}
\Psi\left(u_{n}\right)<r_{n} \tag{18}
\end{equation*}
$$

and

$$
\begin{equation*}
\Phi\left(u_{n}\right)-\inf _{v \in \frac{(\Psi-1}{\left(\Psi^{-\infty}, r_{n}[)\right)_{w}}} \Phi(v)<r_{n}-\Psi\left(u_{n}\right) \tag{19}
\end{equation*}
$$

and in addition,

$$
\begin{equation*}
\liminf _{\|u\| \rightarrow+\infty}(\Psi(u)+\Phi(u))=-\infty \tag{20}
\end{equation*}
$$

then, there exists a sequence $\left(v_{n}\right)_{n}$ of local minima of $\Psi+\Phi$ such that $\Psi\left(v_{n}\right) \rightarrow+\infty$ as $n \rightarrow \infty$.
(c) If there exists a sequence $\left(r_{n}\right)_{n} \subset\left(\inf _{E} \Psi,+\infty\right)$ with $r_{n} \rightarrow \inf _{E} \Psi$ and a sequence $\left(u_{n}\right)_{n} \subset E$ such that for each $n$ the condition (18) and (19) are satisfied, and in addition,
every global minimizer of $\Psi$ is not a local minimizer of $\Phi+\Psi$,
then, there exists a sequence $\left(v_{n}\right)_{n}$ of pairwise distinct local minimizers of $\Phi+\Psi$ such that $\lim _{n \rightarrow \infty} \Psi\left(v_{n}\right)=\inf _{E} \Psi$, and $\left(v_{n}\right)_{n}$ weakly converges to a global minimizer of $\Psi$.

## 3. Basic assumptions and main results

Throughout the paper, we assume that the following assumptions hold true:
Let $\Omega$ be a bounded open subset of $\mathbb{R}^{N}$ with boundary of class $C^{1}$, and let $v_{i}$ be the components of the outer normal unit vector on the boundary $\partial \Omega$.

Assume that $f, g: \Omega \times \mathbb{R} \longmapsto \mathbb{R}$ are Carathéodory functions satisfying,

$$
\begin{equation*}
\sup _{|t| \leq r}|f(x, t)| \in L^{1}(\Omega), \quad \text { and } \quad \sup _{|t| \leq r}|g(x, t)| \in L^{1}(\Omega) \quad \text { for each } \quad r>0 . \tag{22}
\end{equation*}
$$

We set

$$
\begin{equation*}
F(x, t)=\int_{0}^{t} f(x, s) d s, \quad \text { and } \quad G(x, t)=\int_{0}^{t} g(x, s) d s \tag{23}
\end{equation*}
$$

Assume the following assumptions:
(M1) $M_{i}: \mathbb{R}_{+} \longrightarrow \mathbb{R}_{+}$are continuous functions satisfy the condition

$$
M_{i}(t) \geq m \text { for all } t \geq 0
$$

where $m$ is a positive constant.
(M2) There are constants $s^{*}>1$ and $\lambda_{1}, \lambda_{2}>0$ such that

$$
\widehat{M}_{0}(t) \leq \lambda_{1} t^{s^{*}}+\lambda_{2} \text { for all } t \geq 0
$$

where $\widehat{M}_{i}(t)=\int_{0}^{t} M_{i}(s) d s$.
We define, for any $u \in W^{1, \vec{p}(\cdot)}(\Omega)$, the functionals

$$
\begin{align*}
J(u) & =\sum_{i=1}^{N} \widehat{M_{i}}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)+\widehat{M_{0}}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|u|^{p_{0}(x)} d x\right),  \tag{24}\\
\Psi(u) & =J(u)-\int_{\Omega} G(x, u) d x \quad \text { and } \quad \Phi(u)=-\int_{\Omega} F(x, u) d x \tag{25}
\end{align*}
$$

We assume that $G$ satisfies one of the following two conditions:
(G1) There are positive functions $\vartheta_{0}(\cdot), \vartheta_{1}(\cdot) \in L^{1}(\Omega)$ with $\vartheta_{0} \neq 0$ such that

$$
|G(x, t)| \leq \frac{m \vartheta_{0}(x)}{2(N+1)^{\underline{p}-1} \bar{p} C_{0}^{\underline{p}}\left\|\vartheta_{0}\right\|_{L^{1}(\Omega)}}|t| \underline{\underline{p}}+\vartheta_{1}(x)
$$

almost everywhere $x \in \Omega$, and for any $t \geq 0$.
(G2) There are a constants $R>0, \varepsilon \in(0,1)$ and $C_{1}>0$ such that for any $|t| \geq R, \quad \int_{\Omega}|G(x, t)| d x \leq(1-\varepsilon) \widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|t|^{p_{0}(x)} d x\right)+C_{1}$, almost everywhere $x$ in $\Omega$.

Remark 3.1. Assume the hypothesis (M2) and one of the assumptions (G1) and (G2). Then there are positives constants $d_{1}$ and $d_{2}$ such that

$$
\begin{equation*}
\widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|\xi|^{p_{0}(x)} d x\right)-\int_{\Omega} G(x, \xi) d x \leq d_{1}|\xi|^{\bar{p} s^{*}}+d_{2}, \text { for all } \xi \in \mathbb{R} \tag{26}
\end{equation*}
$$

Indeed, on the one hand, assumptions (M2) and (G1) implies

$$
\begin{aligned}
& \widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|\xi|^{p_{0}(x)} d x\right)-\int_{\Omega} G(x, \xi) d x \\
& \leq \lambda_{1}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|\xi|^{p_{0}(x)} d x\right)^{s^{*}}+\lambda_{2}+\frac{m}{2(N+1)^{\underline{p}-1} \bar{p} C_{0}^{\underline{p}}}|\xi| \underline{p}+\left\|\vartheta_{1}\right\|_{L^{1}(\Omega)} \\
& \leq \frac{\lambda_{1}}{\underline{p}^{s^{*}}}|\xi|^{\bar{p} s^{*}}+\frac{m}{2(N+1)^{\underline{p}-1} \bar{p} C_{0}^{\underline{p}}}|\xi| \underline{p}+\left\|\vartheta_{1}\right\|_{L^{1}(\Omega)}+\lambda_{2}, \text { for } \xi \text { large enough } .
\end{aligned}
$$

Then, we have (26) for $d_{1}=\max \left(\frac{\lambda_{1}}{\underline{p}^{p^{*}}}, \frac{m}{2(N+1)^{\underline{p-1}} \bar{p} C_{0}^{p}}\right)$ and $d_{2}=\left\|\vartheta_{1}\right\|_{L^{1}(\Omega)}+\lambda_{2}$. On the other hand, assumptions (M2) and (G2) implies

$$
\begin{aligned}
& \widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|\xi|^{p_{0}(x)} d x\right)-\int_{\Omega} G(x, \xi) d x \\
& \leq \widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|u|^{p_{0}(x)} d x\right)+(1-\varepsilon) \widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|t|^{p_{0}(x)} d x\right)+C_{1} \\
& \leq(2-\varepsilon) \frac{\lambda_{1}}{\underline{p}^{s^{*}}}|\xi|^{\overline{s^{*}}}+C_{1}+\lambda_{2}(2-\varepsilon) \text { for } \xi \text { large enough } .
\end{aligned}
$$

Which ends the proof of (26) with $d_{1}=(2-\varepsilon) \frac{\lambda_{1}}{\underline{p}^{*}}$ and $d_{2}=C_{1}+\lambda_{2}(2-\varepsilon)$.
Let us prove that the functional $\Psi$ is coercive.

Lemma 3.2. Assume that (M1) and (M2) hold and one of the condition (G1) or (G2) is satisfied. Then the functional $\Psi$ is coercive.

Proof. Assuming that the condition (G1) is satisfied, then for $\|u\|_{1, \vec{p}(\cdot)} \geq 1$, we have

$$
\begin{align*}
& \Psi(u)=J(u)-\int_{\Omega} G(x, u) d x \\
& =\sum_{i=1}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)+\widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|u|^{p_{0}(x)} d x\right) \\
& -\int_{\Omega} G(x, u) d x \\
& \geq \sum_{i=1}^{N} \frac{m}{p_{i}^{+}}\left(\left\|\frac{\partial u}{\partial x_{i}}\right\| \|_{p_{i}(\cdot)}^{\underline{p}}-1\right)+\frac{m}{p_{0}^{+}}\left(\|u\|_{p_{0}(x)}^{\frac{p}{p}}-1\right) \\
& -\frac{m}{2(N+1)^{\underline{p}-1} \bar{p} C^{\underline{p}}} \int_{\Omega} \frac{\vartheta_{0}(x)}{\left\|\vartheta_{0}\right\|_{L^{1}(\Omega)}}|u|^{\underline{p}} d x-\int_{\Omega} \vartheta_{1}(x) d x,(\text { by (M1) }) \\
& \geq \frac{m}{\bar{p}(N+1)^{\underline{p}-1}}\|u\|_{1, \vec{p}(\cdot)}^{\frac{p}{p}}-\frac{m}{2(N+1)^{\underline{p}-1} \bar{p} C_{0}^{\underline{p}}}\|u\|_{L^{\infty}(\Omega)}^{\frac{p}{p}} \\
& -\left\|\vartheta_{1}(\cdot)\right\|_{L^{1}(\Omega)}-N-1 \\
& \geq \frac{m}{\bar{p}(N+1)^{\underline{p}-1}}\|u\|_{1, \vec{p}(\cdot)}^{\frac{p}{p}}-\frac{m}{2(N+1)^{\underline{p}-1} \bar{p}}\|u\|_{1, \vec{p}(\cdot)}^{\frac{p}{p}}-C_{2}  \tag{14}\\
& \geq \frac{m}{2 \bar{p}(N+1)^{\underline{p}-1}}\|u\| \frac{p}{1, \vec{p}(\cdot)}-C_{2} . \tag{27}
\end{align*}
$$

Under the condition (G2) we have for $\|u\|_{1, \vec{p}(\cdot)} \geq 1$

$$
\begin{align*}
\Psi(u) & =J(u)-\int_{\Omega} G(x, u) d x \\
& =\sum_{i=1}^{N} \widehat{M_{i}}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)+\widehat{M_{0}}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|u|^{p_{0}(x)} d x\right) \\
& -(1-\varepsilon) \widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|t|^{p_{0}(x)} d x\right)-C_{1} \\
& \geq \sum_{i=1}^{N} \frac{m}{p_{i}^{+}}\left(\left.\left\|\frac{\partial u}{\partial x_{i}}\right\|\right|_{p_{i}(\cdot)} ^{\underline{p}}-1\right)+\frac{m \varepsilon}{p_{0}^{+}}\left(\|u\|_{p_{0}(x)}^{\frac{p}{p}}-1\right)-C_{1} \\
& \geq \frac{\min \{1, \varepsilon\} m}{\bar{p}(N+1)^{\frac{p}{-1}}}\|u\|_{\frac{p}{1, \vec{p}(\cdot)}}-N-m \varepsilon-C_{1} \\
& \geq \frac{\min \{1, \varepsilon\} m}{\bar{p}(N+1)^{\frac{p}{p}-1}}\|u\|_{\frac{p}{1, \vec{p}}(\cdot)}^{\frac{p}{2}}-C_{3} . \tag{28}
\end{align*}
$$

Thanks to (27)-(28) we conclude that $\Psi$ is coercive. Moreover, there exist two positive constants $\alpha$ and $\beta$ such that

$$
\begin{equation*}
\Psi(u) \geq \alpha\|u\|_{1, \vec{p}(\cdot)}^{\frac{p}{}} \quad \text { for all }\|u\|_{1, \vec{p}(\cdot)} \geq \beta \tag{29}
\end{equation*}
$$

Definition 3.3. A measurable function $u \in W^{1, \vec{p}(\cdot)}(\Omega)$ is called a weak solution of the Neumann elliptic problem (8) if

$$
\begin{align*}
& \sum_{i=1}^{N} M_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \int_{\Omega} \sum_{i=1}^{N}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} d x  \tag{30}\\
+ & M_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|u|^{p_{0}(x)} d x\right) \int_{\Omega}|u|^{p_{0}(x)-2} u v d x=\int_{\Omega} f(x, u) v(x) d x+\int_{\Omega} g(x, u) v(x) d x
\end{align*}
$$

for all $v \in W^{1, \vec{p}(\cdot)}(\Omega)$.
Definition 3.4. A function $F(x, t)$ satisfies the condition (S) if for each compact subset $E$ of $\mathbb{R}$, there exists $\xi \in E$ such that
(S)

$$
F(x, \xi)=\sup _{t \in E} F(x, t) \text { for a.e. } x \in \Omega
$$

Now we are in the position to state our first main result.

Theorem 3.5. Let the conditions (M1)-(M2) be satisfied, and let (G1) or (G2) hold. Moreover, let F satisfy the condition (S). Moreover, we suppose that

$$
\begin{equation*}
\liminf _{|\xi| \rightarrow+\infty} \widehat{M}_{0}\left(\int_{\Omega}\left(\frac{1}{p_{0}(x)}|\xi|^{p_{0}(x)}\right) d x\right)-\int_{\Omega}(G(x, \xi)+F(x, \xi)) d x=-\infty \tag{31}
\end{equation*}
$$

and there are positives sequences $\left(a_{n}\right)_{n}$ and $\left(b_{n}\right)_{n}$ such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} b_{n}=+\infty, \quad \lim _{n \rightarrow \infty} \frac{a_{n}^{\bar{p}}}{b_{n}^{\frac{p}{p}}}=0 \tag{32}
\end{equation*}
$$

Finally, we assume that there exist a positive function
$h(\cdot) \in L^{1}(\Omega)$ with $\|h(\cdot)\|_{L^{1}(\Omega)}=1$ and suitable positives constants $d_{1}$ and $d_{2}$ such that for each $n$ we have for almost every $x$ in $\Omega$

$$
\begin{array}{r}
F\left(x, a_{n}\right)+h(x)\left(\alpha\left(\frac{b_{n}}{C_{0}}\right)^{\underline{p} s^{*}}-d_{1} a_{n}^{\bar{p} s^{*}}-d_{2}\right) \geq \sup _{t \in\left[a_{n}, b_{n}\right]} F(x, t), \\
F\left(x,-a_{n}\right)+h(x)\left(\alpha\left(\frac{b_{n}}{C_{0}}\right)^{\underline{p} s^{*}}-d_{1} a_{n}^{\bar{p} s^{*}}-d_{2}\right) \geq \sup _{t \in\left[-b_{n},-a_{n}\right]} F(x, t), \tag{34}
\end{array}
$$

where $\alpha$ is constant of coercivity defined in (29), $d_{1}=\max \left(\frac{\lambda_{1}}{\underline{p}^{* *}}, \frac{m}{2(N+1)^{\underline{p-1}} \bar{p} C_{0}^{\underline{p}}}\right)$ and $d_{2}=\left\|\vartheta_{1}\right\|_{L^{1}(\Omega)}+\lambda_{2}$ if we assume (G1) and $d_{1}=(2-\varepsilon) \frac{\lambda_{1}}{\underline{p}^{*}}$ and $d_{2}=C_{1}+$ $\lambda_{2}(2-\varepsilon)$ if we assume (G2). The last inequalities (33) and (34) are strict on a subset of $\Omega$ with positive measure.
Then there is a sequence $\left(v_{n}\right)_{n}$ of local minima of $\Psi+\Phi$ such that $\lim _{n \rightarrow \infty} \Psi\left(v_{n}\right)=$ $+\infty$. Consequently, the problem (8) admits an unbounded sequence of weak solutions.

Proof. Let $\Psi, \Phi$ be the functionals defined in (23)-(25).
Since $\underline{p}>N$, the embedding $W^{1, \vec{p}(\cdot)}(\Omega) \hookrightarrow C^{0}(\bar{\Omega})$ is continuous and compact. We can see that $J, \Phi, \Psi \in \mathcal{C}^{1}\left(W^{1, \vec{p} \cdot \cdot)}(\Omega), \mathbb{R}\right)$ (see [12], [29]) with the derivative given by

$$
\begin{align*}
\left\langle\Psi^{\prime}(u), v\right\rangle & =\sum_{i=1}^{N} M_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \int_{\Omega} \sum_{i=1}^{N}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} d x \\
& +M_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|u|^{p_{0}(x)} d x\right) \int_{\Omega}|u|^{p_{0}(x)-2} u v d x-\int_{\Omega} g(x, u) v(x) d x \tag{35}
\end{align*}
$$

and

$$
\begin{equation*}
\left\langle\Phi^{\prime}(u), v\right\rangle=-\int_{\Omega} f(x, u) v(x) d x, \quad \text { for all }(u, v) \in W^{1, \vec{p}(\cdot)}(\Omega) \tag{36}
\end{equation*}
$$

Then $u \in W^{1, \vec{p}(\cdot)}(\Omega)$ is a weak solution of (8) if and only if $u$ is a critical point of $\Psi+\Phi$.

Let us prove that the functionals $\Psi$ and $\Phi$ are sequentially weakly lower semi-continuous.

For $i=0, \ldots, N$ and any $u \in W^{1, \vec{p}(\cdot)}(\Omega)$ define $J_{i}, H: W^{1, \vec{p}(\cdot)}(\Omega) \longrightarrow \mathbb{R}$ by

$$
\begin{gathered}
J_{i}=\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x, \quad \text { where } \quad \frac{\partial u}{\partial x_{0}}=u \\
H(u)=-\int_{\Omega} G(x, u) d x
\end{gathered}
$$

Let $\left(u_{n}\right)_{n}$ be a sequence such that $u_{n} \rightharpoonup u$ in $W^{1, \vec{p}(\cdot)}(\Omega)$. Since $J_{i}$ is convex, for any $n$ we have

$$
J_{i}(u) \leq J_{i}\left(u_{n}\right)+\left\langle J_{i}^{\prime}(u), u-u_{n}\right\rangle .
$$

Passing to the limit in the above inequality with $n \rightarrow \infty$, we see that $J_{i}$ is sequentially weakly lower semi-continuous. Then we have

$$
\begin{equation*}
\int_{\Omega} \sum_{i=1}^{N} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x \leq \liminf _{n \rightarrow+\infty} \int_{\Omega} \sum_{i=0}^{N} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x \tag{37}
\end{equation*}
$$

From (37) and since $\widehat{M_{i}}$ is continuous and monotone, we have

$$
\begin{align*}
\liminf _{n \rightarrow+\infty} J\left(u_{n}\right) & =\liminf _{n \rightarrow+\infty} \sum_{i=0}^{N} \widehat{M_{i}}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \\
& \geq \sum_{i=0}^{N} \widehat{M_{i}}\left(\liminf _{n \rightarrow+\infty} \int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \\
& \geq \sum_{i=0}^{N} \widehat{M}_{i}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right) \\
& \geq J(u), \tag{38}
\end{align*}
$$

namely, $J$ is sequentially weakly lower semi-continuous.
On the other hand, by Remark 2.2 up to a sub-sequence we have $u_{n} \longrightarrow u$ in $C^{0}(\bar{\Omega})$. Hence,

$$
\begin{aligned}
u_{n} & \longrightarrow u \text { uniformelly in } \Omega \\
k & =\sup _{n \in \mathbb{N}}\left\|u_{n}\right\|_{L^{\infty}(\Omega)}<+\infty
\end{aligned}
$$

Therefore, $G\left(x, u_{n}(x)\right) \longrightarrow G(x, u(x))$ almost every $x$ in $\Omega$ and $\left|G\left(x, u_{n}(x)\right)\right| \leq$ $k \sup |g(x, s)|$. Note that $\sup |g(x, s)| \in L^{1}(\Omega)$ by (22). Thus, the dominated $|s| \leq k$
convergence theorem implies that $\lim _{n \rightarrow \infty} H\left(u_{n}\right)=H(u)$. So, the functional $H$ is sequentially weakly continuous on $W^{1, \vec{p}(\cdot)}(\Omega)$, and hence, being $\Psi=J-H, \Psi$ is sequentially weakly lower semi-continuous. In the same way (as in the case of the mapping $H$ ) we can show that $\Phi$ is sequentially weakly continuous.

For $\rho>\inf _{u \in W^{1, \vec{p}(\cdot)}(\Omega)} \Psi(u)$, we define

$$
\begin{equation*}
\mathbf{K}(\rho)=\inf \left\{\tau>0 \text { suh that } \Phi^{-1}(]-\infty, \rho[) \subset \overline{B(0, \tau)}\right\} \tag{39}
\end{equation*}
$$

where $B(0, \tau)=\left\{u \in W^{1, \vec{p}(\cdot)}(\Omega):\|u\|_{1, \vec{p}(\cdot)}<\tau\right\}$ and $\overline{B(0, \tau)}$ denotes the closure of $B(0, \tau)$ in $W^{1, \vec{p}(\cdot)}(\Omega)$ for the norm topology.

Owing to the fact that $\Psi$ is coercive, we have $0<\mathbf{K}(\rho)<+\infty$ for each $\rho>\inf _{u \in W^{1, \vec{p} \cdot()}(\Omega)} \Psi(u)$. In view of (29), we deduce that

$$
\text { if } \Psi(u)<\alpha\|u\|_{1, \vec{p}(\cdot)}^{\frac{p}{p}} \text {, then }\|u\|_{1, \vec{p}(\cdot)}<\beta
$$

Thanks to (39), one has $\Psi^{-1}(]-\infty, \rho[) \subset \overline{B(0, \mathbf{K}(\rho))}$ and so $\overline{\left(\Psi^{-1}(]-\infty, \rho[)\right)_{w}} \subset \overline{B(0, \mathbf{K}(\rho))}$.
Using (14), we get $\|u\|_{L^{\infty}(\Omega)} \leq C_{0}\|u\|_{1, \vec{p}(\cdot)}$. Then,

$$
\overline{B(0, \mathbf{K}(\rho))} \subset\left\{u \in C(\bar{\Omega}):\|u\|_{L^{\infty}(\Omega)} \leq C_{0} \mathbf{K}(\rho)\right\}
$$

which yields

$$
\begin{equation*}
\inf _{v \in \frac{\left(\Psi^{-1}(]-\infty, \rho[)\right)_{w}}{}} \Phi(v) \geq \inf _{\|v\|_{1, \vec{p} \cdot()} \leq \mathbf{K}(\rho)} \Phi(v) \geq \inf _{\|v\|_{L^{\infty}(\Omega)} \leq C_{0} \mathbf{K}(\rho)} \Phi(v) . \tag{40}
\end{equation*}
$$

Let $\tau \geq \alpha \beta^{p s^{*}}$ and $u \in W^{1, \vec{p}(\cdot)}(\Omega)$ be such that $\Psi(u)<\tau$. When $\|u\|_{1, \vec{p}(\cdot)} \geq \beta$, by (29), one has

$$
\tau>\Psi(u) \geq \alpha\|u\|_{1, \vec{p} \cdot(\cdot)}^{\frac{p}{p}}
$$

which implies that $\|u\|_{1, \vec{p}(\cdot)} \leq\left(\frac{\tau}{\alpha}\right)^{\frac{1}{p^{s^{*}}}}$. When $\|u\|_{1, \vec{p}(\cdot)}<\beta$, it is clear that $\|u\|_{1, \vec{p}(\cdot)} \leq\left(\frac{\tau}{\alpha}\right)^{\frac{1}{p s^{*}}}$.

By the definition of $\mathbf{K}(\tau)$, we have

$$
\begin{equation*}
\mathbf{K}(\tau) \leq\left(\frac{\tau}{\alpha}\right)^{\frac{1}{\underline{p}^{s^{*}}}} \tag{41}
\end{equation*}
$$

Since $F(x, \cdot)$ satisfies condition (S), for each $n$, there exists $\xi_{n} \in\left[-a_{n}, a_{n}\right]$ such that

$$
\begin{equation*}
F\left(x, \xi_{n}\right)=\sup _{t \in\left[-a_{n}, a_{n}\right]} F(x, t) \text { a.e. in } \Omega . \tag{42}
\end{equation*}
$$

Now, in order to satisfy (b) of Theorem 2.3, take as $u_{n}$ the constant function whose value is $\xi_{n}$ and $\rho_{n}=\alpha\left(\frac{b_{n}}{C_{0}}\right)^{\underline{p} s^{*}}$, then $\lim _{n \rightarrow \infty} \rho_{n} \rightarrow+\infty$, and thanks to (41) we obtain

$$
\begin{equation*}
\mathbf{K}\left(\rho_{n}\right) \leq \frac{b_{n}}{C_{0}} \quad \text { then } \quad C_{0} \mathbf{K}\left(\rho_{n}\right) \leq b_{n} \tag{43}
\end{equation*}
$$

By (26), one has

$$
\begin{aligned}
e_{n} & =\widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}\left|\xi_{n}\right|^{p_{0}(x)} d x\right)-\int_{\Omega} G\left(x, \xi_{n}\right) d x \\
& \leq d_{1}\left|\xi_{n}\right|^{\mid \overline{p s}}+d_{2} \leq d_{1}\left|a_{n}\right|^{\mid \bar{s} s^{*}}+d_{2}
\end{aligned}
$$

It follows from (32) that for $n$ large enough,

$$
d_{1}\left|a_{n}\right|^{\bar{p} s^{*}}+d_{2}<\alpha\left(\frac{b_{n}}{C_{0}}\right)^{\underline{p s^{*}}}=\rho_{n}
$$

and consequently $e_{n}<\rho_{n}$, that is (18) holds. Without loss of generality, we may assume that (18) holds for all $n$.

From (33)-(34) and (42), we obtain

$$
\begin{equation*}
F\left(x, \xi_{n}\right)+h(x)\left(\rho_{n}-e_{n}\right) \geq \sup _{|t| \leq b_{n}} F(x, t) \text { a.e. in } \Omega \tag{44}
\end{equation*}
$$

and the inequality (44) is strict on a subset of $\Omega$ with positive measure. Using (43) and (44), we obtain (19) and (20) follows directly from (31).

Therefore, all hypotheses of Theorem 2.3 (b) are satisfied, then the proof of the Theorem 3.5 is concluded.

Our second result is the following theorem.
Theorem 3.6. Assume that (M1)-(M2) hold. Suppose that

$$
\begin{equation*}
G(x, t) \leq 0 \text { for } t \in \mathbb{R} \text { and a.e. } x \in \Omega \tag{45}
\end{equation*}
$$

and that there exist two positive constants $\delta$ and $\varepsilon$ such that

$$
\begin{equation*}
-G(x, t) \leq \delta|t| \underline{p} \text { for }|t| \leq \varepsilon \text { and a.e. } x \in \Omega \tag{46}
\end{equation*}
$$

Moreover, let the functional $F$ satisfy the condition $(\mathbf{S})$ and

$$
\begin{equation*}
\limsup _{|\xi| \rightarrow 0} \frac{\int_{\Omega} F(x, \xi) d x+\int_{\Omega} G(x, \xi) d x}{|\xi|^{\underline{p}}}>\widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)} d x\right) \tag{47}
\end{equation*}
$$

Suppose that $\left(a_{n}\right)_{n}$ and $\left(b_{n}\right)_{n}$ be two positives sequences such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} b_{n}=0 \quad \text { and } \quad \lim _{n \rightarrow \infty} \frac{a_{n}^{\frac{p}{n}}}{b_{n}^{\bar{p}}}=0 \tag{48}
\end{equation*}
$$

and there exists a positive function $h \in L^{1}(\Omega)$ with $\|h\|_{L^{1}(\Omega)}=1$ such that for each $n$ and almost every $x$ in $\Omega$, we have

$$
\begin{array}{r}
F\left(x, a_{n}\right)+h(x)\left(d_{3}\left(\frac{b_{n}}{C_{0}}\right)^{\bar{p} s^{*}}-d_{4} a_{n}^{\underline{p} s^{*}}-\lambda_{2}\right) \geq \sup _{t \in\left[a_{n}, b_{n}\right]} F(x, t), \\
F\left(x,-a_{n}\right)+h(x)\left(d_{3}\left(\frac{b_{n}}{C_{0}}\right)^{\bar{p} s^{*}}-d_{4} a_{n}^{p s^{*}}-\lambda_{2}\right) \geq \sup _{t \in\left[-b_{n},-a_{n}\right]} F(x, t), \tag{50}
\end{array}
$$

with $d_{3}=\frac{m}{\bar{p}(N+1)^{\bar{p}-1}}$ and $d_{4}=\max \left(\frac{\lambda_{1}}{\underline{p}^{p^{*}}}, \delta|\Omega|\right)$, the inequalities (49) and (50) are strict on a subset of $\Omega$ with positive measure.

Then there exists a sequence $\left(v_{n}\right)_{n}$ of pairwise distinct local minima of $\Psi+\Phi$ such that $v_{n} \rightarrow 0$ in $W^{1, \vec{p}(\cdot)}(\Omega)$. Consequently, the problem (8) admits a sequence of non-zero weak solutions which strongly converges to 0 in $W^{1, \vec{p}(\cdot)}(\Omega)$.

Proof. Let us verify all the hypotheses of Theorem 2.3 point (c). Using (45), for $\|u\|_{1, \vec{p}(\cdot)} \leq 1$ we have

$$
\begin{aligned}
\Psi(u) & =J(u)-\int_{\Omega} G(x, u) d x \\
& =\sum_{i=1}^{N} \widehat{M_{i}}\left(\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)+\widehat{M_{0}}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|u|^{p_{0}(x)} d x\right) \\
& \geq d_{3}\|u\|_{1, \vec{p}(\cdot)}^{\bar{p}}-C_{4}
\end{aligned}
$$

with $d_{3}=\frac{m}{\bar{p}(N+1)^{\bar{p}-1}}$. Then, $\Psi$ is coercive, $\inf _{W^{1, p(\cdot)}(\Omega, w)} \Psi=\Psi(0)=0$ and 0 is the unique global minimizer of $\Psi$. Thanks to (47) we have

$$
\begin{aligned}
& \limsup _{|\xi| \rightarrow 0}\{\Psi(\xi)+\Phi(\xi)\} \\
& \quad=\limsup _{|\xi| \rightarrow 0}\left\{\widehat{M}_{0}\left(\int_{\Omega} \frac{|\xi|^{p_{0}(x)}}{p_{0}(x)} d x\right)-\int_{\Omega} G(x, \xi) d x-\int_{\Omega} F(x, \xi) d x\right\} \\
& \quad \leq \limsup _{|\xi| \rightarrow 0}\left\{\widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)} d x\right)|\xi| \underline{p}-\int_{\Omega} G(x, \xi) d x-\int_{\Omega} F(x, \xi) d x\right\}<0
\end{aligned}
$$

that is, 0 is not a local minimizer of $\Psi+\Phi$; so (21) is satisfied.
For $r>0$ sufficiently small, the condition $\Psi(u)<r$ implies that $\|u\|_{1, \vec{p}(\cdot)}<$ $\left(\frac{r}{d_{3}}\right)^{\frac{1}{\overline{\beta^{*}}}}$, this shows that

$$
\mathbf{K}(r) \leq\left(\frac{r}{d_{3}}\right)^{\frac{1}{\overline{p s} s^{*}}}
$$

Now put $\rho_{n}=d_{3}\left(\frac{b_{n}}{C_{0}}\right)^{\bar{p} s^{*}}$ and take $u_{0}$ and $u_{n}$ as constants values functions $\xi_{0}$ and $\xi_{n}$ respectively in Theorem 2.3. Then

$$
\begin{equation*}
C_{0} \mathbf{K}\left(\rho_{n}\right) \leq b_{n} \tag{51}
\end{equation*}
$$

By (46) and (M2), there exists a sequence $\left(\xi_{n}\right)_{n} \subset \mathbb{R}$ with $\xi_{n} \in\left[-a_{n}, a_{n}\right]$ such that for each $a_{n}$ sufficiently small,

$$
\begin{align*}
e_{n} & =\widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p(x)}\left|\xi_{n}\right|^{p(x)} d x\right)-\int_{\Omega} G\left(x, \xi_{n}\right) d x \\
& \leq \widehat{M}_{0}\left(\int_{\Omega} \frac{\left|\xi_{n}\right|^{p_{0}(x)}}{p_{0}(x)} d x\right)+\delta|\Omega|\left|\xi_{n}\right|^{\underline{p}} \\
& \leq \frac{\lambda_{1}}{\underline{p}^{s^{*}}}\left|\xi_{n}\right|^{\underline{p}^{*}}+\delta|\Omega|\left|\xi_{n}\right|^{\underline{p}}+\lambda_{2} \\
& \leq d_{4}\left|\xi_{n}\right|^{\underline{p}}+\lambda_{2} \\
& \leq d_{4}\left|a_{n}\right|^{\underline{p}}+\lambda_{2} \tag{52}
\end{align*}
$$

where $d_{4}=\max \left(\frac{\lambda_{1}}{p^{s^{*}}}, \delta|\Omega|\right)$.
It follows from (48) that for $n$ large enough,

$$
d_{4}\left|a_{n}\right| \underline{s}^{s^{*}}<d_{3}\left(\frac{b_{n}}{C_{0}}\right)^{\bar{p} s^{*}}=\rho_{n}+\lambda_{2}
$$

Then (18) is obtained.
Noting that $F(x, \cdot)$ satisfies condition (S), for each $n$, there exists $\xi_{n} \in\left[-a_{n}, a_{n}\right]$ such that

$$
\begin{equation*}
F\left(x, \xi_{n}\right)=\sup _{t \in\left[-a_{n}, a_{n}\right]} F(x, t) \quad \text { a.e. in } \Omega . \tag{53}
\end{equation*}
$$

Then, thanks to (49) and (50) we can obtain that

$$
\begin{equation*}
F\left(x, \xi_{n}\right)+h(x)\left(\rho_{n}-e_{n}-\lambda_{2}\right) \geq \sup _{|t| \leq b_{n}} F(x, t) \text { a.e. in } \Omega \tag{54}
\end{equation*}
$$

and the inequality (54) is strict on a subset of $\Omega$ with positive measure. Thanks to (51) and (54) we obtain (19). Therefore, the hypotheses of Theorem 2.3 (c) are satisfied.

Consequently, there exists a sequence $\left(v_{n}\right)_{n}$ of pairwise distinct local minima of $\Psi+\Phi$ such that $\Psi\left(v_{n}\right) \rightarrow 0$, which implies $\left\|v_{n}\right\|_{1, \vec{p}(\cdot)} \rightarrow 0$, which complete the proof.

Now we give some remarks and some examples which motivate our results.
Remark 3.7. The definition of our framework requires only the measurability of $p(\cdot)$ and we do not use Modular-Poincaré inequality which require the log-Hölder continuity of the exponent $p(\cdot)$, while the norm Poincaré inequality requires only the continuity of $p(\cdot)$, for more details we refer to [3], [18].

Remark 3.8. Observe that our results require the condition (S) on $F$ (for more details on the condition ( $\mathbf{S}$ ) see [21], Remark 2.5), and on the function $G$ we need one of the assumptions (G1) which is a growth condition on $G$ and (G2) which gives a relation between $\widehat{M}_{0}$ and the integral of $G$.

Now, we present some examples to illustrate our results.
For $i \in[[0, N]]$ we set $M_{i}(t)=(1+t)^{\theta_{i}-1}$ where $\theta_{i}>1$ the we have $M_{i}(t) \geq 1$ for all $t \geq 0$ and by using following classical result

$$
(a+b)^{\theta} \leq 2^{\theta-1}\left(a^{\theta}+b^{\theta}\right), \quad \text { for all } a, b \geq 0 \text { and } \theta \geq 1
$$

we get

$$
\widehat{M}_{0}(t)=\frac{(1+t)^{\theta_{0}}}{\theta_{0}} \leq 2^{\theta_{0}-1}\left(1+t^{\theta_{0}}\right)
$$

Which means that (M1) and (M2) hold with $m=1, s^{*}=\theta_{0}$ and $\lambda_{1}=\lambda_{2}=2^{\theta_{0}-1}$.
Proposition 3.9. Let $g(x, t)=\frac{\underline{p} \vartheta_{0}(x)}{2(N+1)^{\underline{p}-1} \bar{p} C_{0}^{p}}|t|^{\underline{p}-1}$ where $\vartheta_{0} \in L^{1}(\Omega)$ is positive function with $\left\|\vartheta_{0}\right\|_{L^{1}(\Omega)}=1$, and let $f(x, t) \equiv \alpha(x) f_{1}(t)$, with $\alpha(\cdot) \in L^{1}(\Omega)$ be a positive function such that $\|\alpha(\cdot)\|_{L^{1}(\Omega)}=1$ and $f_{1}(\cdot)$ a continuous function with $f_{1}(t)=F_{1}^{\prime}(t)$ and $F_{1}(-t)=F_{1}(t)$. Then, the following nonlinear perturbed Kirchhoff problem

$$
\left\{\begin{array}{l}
-\sum_{i=1}^{N}\left(1+\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)^{\theta_{i}-1} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}}\right)  \tag{55}\\
+\left(1+\int_{\Omega} \frac{1}{p_{0}(x)}|u|^{p_{0}(x)}\right)^{\theta_{0}-1}|u|^{p_{0}(x)-2} u \\
=\alpha_{1}(x) f_{1}(t)+\frac{p \underline{\vartheta_{0}(x)}}{2(N+1)^{\underline{p}-1} \bar{p} c_{0}^{p}}|u|^{\underline{p-1}} \text { in } \Omega \\
\sum_{i=1}^{N}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}} v_{i}=0 \quad \text { on } \quad \partial \Omega
\end{array}\right.
$$

admits a sequence of weak solutions $\left(u_{n}\right)_{n}$ in $W^{1, \vec{p}(x)}(\Omega)$ such that

$$
\lim _{n \rightarrow \infty}\left\|u_{n}\right\|_{1, \vec{p}(x)}=\infty
$$

Proof. It is clear that $\Psi$, is coercive. We have $F(x, t)=\alpha(x) F_{1}(t)$, Choose two positive sequences $\left(a_{n}\right)_{n}$ and $\left(b_{n}\right)_{n}$ such that $a_{1} \geq 1, b_{n}^{p}=2 n^{2} a_{n}^{\bar{p}}$ and $a_{n+1}>b_{n}$ for every $n$. Define $F_{1}\left(a_{n}\right)=a_{n}^{\theta_{0} \bar{p}+1}$ and $F_{1}\left(b_{n}\right)$ such that

$$
\begin{equation*}
F_{1}\left(a_{n}\right)<F_{1}\left(b_{n}\right)<\left(\frac{1}{2 \bar{p}(N+1)^{\underline{p}-1}}\left(\frac{b_{n}}{C_{0}}\right)^{\underline{p} \theta_{0}}-d_{1}\left|a_{n}\right|^{\bar{p} \theta_{0}}-d_{2}\right)+F_{1}\left(a_{n}\right) \tag{56}
\end{equation*}
$$

where $d_{1}=\max \left(\frac{2^{\theta_{0}-1}}{\underline{p}^{\theta_{0}}}, \frac{1}{2(N+1)^{\underline{p}-1} \bar{p} c_{0}^{\underline{p}}}\right)$ and $d_{2}=\left\|\vartheta_{1}\right\|_{L^{1}(\Omega)}+2^{\theta_{0}-1}$.

$$
\text { Put } \rho_{n}=\frac{1}{2 \bar{p}(N+1)^{\underline{p}-1}}\left(\frac{b_{n}}{C_{0}}\right)^{\underline{p}^{\theta_{0}}} \text { and } \xi_{n}=a_{n}
$$

Since

$$
\begin{aligned}
& \widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}\left|a_{n}\right|^{p_{0}(x)} d x\right)-\int_{\Omega} F\left(x, a_{n}\right) d x \\
\leq & d_{1}\left|a_{n}\right|^{\bar{p} \theta_{0}}+d_{2}-\|\alpha(\cdot)\|_{L^{1}(\Omega)} a_{n}^{\bar{p}^{+}+1} \rightarrow-\infty
\end{aligned}
$$

as $n \rightarrow \infty$, then the conditions (31) - (32) holds true. Taking $w_{1}(x)=\alpha(x)$, and in view of (56) we can obtain the conditions (33) - (34).
The hypotheses of Theorem 3.5 are satisfied, then the problem (55) admits a sequence of weak solutions $\left(u_{n}\right)_{n}$ in $W^{1, \vec{p}(x)}(\Omega)$ such that $\lim _{n \rightarrow \infty}\left\|u_{n}\right\|_{1, \vec{p}(x)}=\infty$.

Remark 3.10. Observe that the function

$$
G(x, t)=\frac{1}{2|\Omega|} \widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)}|t|^{p_{0}(x)} d x\right)+C_{1} \gamma(x)
$$

where $\|\gamma\|_{L^{1}(\Omega)}=1$ satisfies the assumption (G2). Then, we get the same result as the previous example with a function satisfies (G2).

Proposition 3.11. Let $g(x, t)=-\underline{p} \mu(x)|t|^{\underline{p}-1}$ where $\mu(\cdot)$ is a positive function such that $\int_{\Omega} \mu(x) d x=1$ and let $f$ be the function defined in the previous proposition, and consider two positives sequences $\left(a_{n}\right)_{n}$ and $\left(b_{n}\right)_{n}$ such that $a_{n}^{p \theta_{0}}=\frac{1}{n} b_{n}^{\bar{p} s^{*}}$ and $b_{n+1}<a_{n}$, with $F_{1}(0)=0, F_{1}\left(a_{n}\right)=a_{n}^{p \theta_{0}+1}$ and

$$
\begin{equation*}
F_{1}\left(a_{n}\right)<F_{1}\left(b_{n}\right)<\left(\left.\frac{1}{\bar{p}(N+1)^{\bar{p}-1}}\left(\frac{b_{n}}{C_{0}}\right)^{\bar{p} \theta_{0}}-d_{3} \right\rvert\, a_{n} \underline{\underline{p}}^{\theta_{0}}\right)+F_{1}\left(a_{n}\right) \tag{57}
\end{equation*}
$$

Then, the following perturbed Kirchhoff problem

$$
\left\{\begin{array}{l}
-\sum_{i=1}^{N}\left(1+\int_{\Omega} \frac{1}{p_{i}(x)}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)} d x\right)^{\theta_{i}-1} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}}\right) \\
+\left(1+\int_{\Omega} \frac{1}{p_{0}(x)}|u|^{p_{0}(x)}\right)^{\theta_{0}-1}|u|^{p_{0}(x)-2} u  \tag{58}\\
=\alpha_{1}(x) f_{1}(t)-\mu(x) \underline{p}|u|^{\underline{p}-1} \text { in } \Omega \\
\sum_{i=1}^{N}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}(x)-2} \frac{\partial u}{\partial x_{i}} v_{i}=0 \quad \text { on } \quad \partial \Omega
\end{array}\right.
$$

admits a sequence of weak solutions $\left(u_{n}\right)_{n}$ in $W^{1, \vec{p}(x)}(\Omega)$ such that

$$
\lim _{n \rightarrow \infty}\left\|u_{n}\right\|_{1, \vec{p}(x)}=0
$$

Proof. Set $\rho_{n}=\frac{1}{\bar{p}(N+1)^{\bar{p}-1}}\left(\frac{b_{n}}{C_{0}}\right)^{\underline{p} \theta_{0}}$ and $\xi_{n}=a_{n}$, we have

$$
\begin{aligned}
& \int_{\Omega} F\left(x, a_{n}\right) d x+\int_{\Omega} G\left(x, a_{n}\right) d x-\widehat{M}_{0}\left(\int_{\Omega} \frac{1}{p_{0}(x)} d x\right)\left|a_{n}\right|^{\underline{p}} \\
& >\left|a_{n}\right|^{\underline{p} \theta_{0}+1}-\left|a_{n}\right|^{\underline{p}}-\frac{\lambda|\Omega|}{\bar{p}}\left|a_{n}\right|^{\underline{p}} \longrightarrow 0 \text { as } n \rightarrow \infty,
\end{aligned}
$$

therefore (45)-(47) hold true. Using (57) we obtain the conditions (49)-(50).
Thus all the assumptions of Theorem 3.6 are satisfied, which completes the proof.

## Acknowledgments

The author would like to thank Professor B. Ricceri and the referees for their remarks and suggestions.

## REFERENCES

[1] E. Acerbi, G. Mingione, Regularity results for a class of functionals with nonstandard growth, Arch. Ration. Mech. Anal. 156 (2001), no. 2, 121-140.
[2] A. Ahmed, T.Ahmedatt, H. Hjiaj and A. Touzani,Existence of infinitly many weak solutions for some quasi-linear $\vec{p}(\cdot)$-elliptic Neumann problems, Math. Bohem. 142 (2017), no. 3, 243-262.
[3] A. Ahmed, H. Hjiaj and A. Touzani, Existence of infinitely many weak solutions for a Neumann elliptic equations involving the $\vec{p}(\cdot)$-Laplacian operator, Rend. Circ. Mat. Palermo (2). 64 (2015), no. 3, 459-473.
[4] A. Ahmed and M. S. B. Elemine Vall, Three weak solutions for a Neumann elliptic equations involving the $\vec{p}(x)$-Laplacian operator, Nonauton. Dyn. Syst. 7 (2020), 224-236.
[5] A. Ahmed, M. S. B. Elemine Vall, A. Touzani, A. Benkirane, Existence of Infinitely Many Solutions the Neumann Problem for quasi-linear Elliptic Systems Involving the $p(\cdot)$ and $q(\cdot)$-Laplacian, Nonlinear Stud. 24 (2017), no. 3, 687-698.
[6] A. Ahmed, M. S. B. Elemine Vall, A. Touzani, A. Benkirane, Infinitely Many Solutions to the Neumann Problem for Elliptic systems in Anisotropic Variable Exponent Sobolev Spaces, Marrocain journal of pure and applied analysis, 3 (2017), no. 1, 70-82.
[7] C. O. Alves, A. El Hamidi, Existence of solution for an anisotropic equation with critical exponent, Differential Integral Equations 21 (2008), no. 1-2, 25-40.
[8] G. Anello, G. Cordaro, An existence theorem for the Neumann problem involving the p-Laplacian, J. Convex Anal. 10 (2003), no. 1, 185-198.
[9] S. N. Antontsev, J. F. Rodrigues, On stationary thermorheological viscous flows, Ann. Univ. Ferrara Sez. VII Sci. Mat. 52 (2006), 19-36.
[10] S. Antontsev, S. Shmarev, Elliptic equations with anisotropic nonlinearity and nonstandard growth conditions, Handbook of Differential Equations: Stationary Partial Differential Equations 3 (2006), 1-100.
[11] A. Arosio, S. Panizzi, On the well-posedness of the Kirchhoff string, Trans. Amer. Math. Soc. 348 (1996) 305-330.
[12] M. Avci, R. A. Mashiyev and B. Cekic, Solutions of an anisotropic nonlocal problem involving variable exponent, Adv. Nonlinear Anal. 2 (2013), 325-338.
[13] M. Bendahmane, M. Chrif and S. El Manouni, An Approximation Result in Generalized Anisotropic Sobolev Spaces and Application, Z. Anal. Anwend. 30 (2011), no. 3, 341-353.
[14] M. Bohner, G. Caristi, F. Gharehgazlouei, and S. Heidarkhani, Existence and Multiplicity of Weak Solutions for a Neumann Elliptic Problem with $\vec{p}(\cdot)$-Laplacian, Nonauton. Dyn. Syst. 7 (2020), no. 1, 53-64.
[15] M. M. Boureanu, A. Matei and M. Sofonea, Nonlinear problems with $p(\cdot)$-growth conditions and applications to antiplane contact models, Adv. Nonlinear Stud. 14 (2014), no. 2, 295-313.
[16] Y. Chen, S. Levine and R. Rao, Variable exponent linear growth functionals in image restoration, SIAM J. Appl. Math. 66 (2006), no. 4, 1383-1406.
[17] F. J. S. A. Corrêa, R. G. Nascimento, On a nonlocal elliptic system of p-Kirchhoff type under Neumann boundary condition, Math. Comput. Modelling 49 (2009) 598-604.
[18] L. Diening, P. Harjulehto, P. Hästö and M. Růžička, Lebesgue and Sobolev spaces with variable exponents, Lecture Notes in Math. 2017 (2011).
[19] M. S. B. Elemine Vall, A. Ahmed, Multiplicity of solutions for a class of Neumann elliptic systems in anisotropic Sobolev spaces with variable exponent, Adv. Oper. Theory 4 (2019), no. 2, 497-513.
[20] X. L. Fan, Anisotropic variable exponent Sobolev spaces and $\vec{p}(\cdot)$-Laplacian equations, Complex Var. Elliptic Equ. 56 (2011), no. 7-9, 623-642.
[21] X. L. Fan, C. Ji, Existence of infinitely many solutions for a Neumann problem involving the $p(x)$-Laplacian, C. J. Math. Anal. Appl. 334 (2007) 248-260.
[22] I. Fragalà, F. Gazzola, and B. Kawohl, Existence and nonexistence results for anisotropic quasi-linear equations, Ann. Inst. H. Poincaré, Analyse Non Linéaire 21 (2004), 715-734.
[23] J. R. Graef, S. Heidarkhani, L. Kong, A variational approach to a Kirchhoff-type problem involving two parameters, Results Math. 63 (2013) 877-889.
[24] P. Harjulehto and P. Hästö, Sobolev Inequalities for Variable Exponents Attaining the Values 1 and n, Publ. Mat., 52 (2008), no. 2, 347-363.
[25] C. Ji, An eigenvalue of an anisotropic quasi-linear elliptic equation with variable exponent and Neumann boundary condition, Nonlinear Anal. 71 (2009) 45074514.
[26] G. Kirchhoff, Mechanik, Teubner, Leipzig, Germany, (1883).
[27] B. Kone, S. Ouaro, S. Traore, Weak solutions for anisotropic nonlinear elliptic equations with variable exponents, Electron. J. Differential Equations. 144 (2009), 1-11.
[28] M. Mihăilescu, P. Pucci and V. Rădulescu, Eigenvalue problems for anisotropic quasilinear elliptic equations with variable exponent, J. Math. Anal. Appl. 340 (2008), no. 1, 687-698.
[29] A. Ourraoui, Multiplicity Of Solutions For $\vec{p}(\cdot)$-Laplacian Elliptic Kirchhoff Type Equations, Appl. Math. E-Notes 20 (2020), 124-132.
[30] J. Rákosník, Some remarks to anisotropic Sobolev spaces I, Beiträge zur Analysis 13 (1979), 55-68.
[31] J. Rákosník, Some remarks to anisotropic Sobolev spaces II, Beiträge zur Analysis 15 (1981), 127-140.
[32] B. Ricceri, A general variational principle and some of its applications, J. Comput. Appl. Math. 113 (2000), no. 1-2, 401-410.
[33] B. Ricceri, Infinitely many solutions of the Neumann problem for elliptic equations involving the p-Laplacian, Bull. Lond. Math. Soc. 33 (2001), no. 3, 331-340.
[34] B. Ricceri, On an elliptic Kirchhoff-type problem depending on two parameters, J. Global Optim. 46 (2010), no. 4, 543-549.
[35] M. Růžička, Electrorheological fluids: modeling and mathematical theory, Lecture Notes in Math. 1748, (2000).
[36] S. M. Shahruz, S. A. Parasurama, Suppression of vibration in the axially moving Kirchhoff string by boundary control, J. Sound Vib. 214 (1998), no. 3, 567-575.
[37] Y. T. Shen, S. S. Yan, The Variational Methods for quasi-linear Elliptic Equations, South China University of Technology Press, Guang Zhou (1995) (in Chinese).
[38] R. Stanway, J. L. Sproston, A. K. El-Wahed, Applications of electro-rheological fluids in vibration control: a survey, Smart Mater. Struct. 5 (1996), no. 4, 464-482.
[39] V. V. Zhikov, Averaging of functionals in the calculus of variations and elasticity, Math. USSR Izv. 29 (1987), no. 1, 33-66.
[40] V. V. Zhikov, On some variational problems, Russ. J. Math. Phys. 5 (1997), no. 1, 105-116.

> A. AHMED
> University of Sidi Mohamed Ibn Abdellah,
> Faculty of Sciences Dhar El Mahraz,
> Laboratory LAMA, Department of Mathematics,
> B.P. 1796 Atlas Fez, Morocco.
> e-mail: ahmedmath2001@gmail .com
M.S.B. ELEMINE VALL

University of Nouakchott, Professional University Institute
Department of Mathematics, Nouakchott, Mauritania.
e-mail: saad2012bouh@gmail.com

