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BIVARIATE EXPONENTIAL INTEGRALS AND EDGE-BICOLORED GRAPHS

M. BORINSKY - C. MERONI - M. WIESMANN

We show that specific exponential bivariate integrals serve as generating functions of labeled edge-bicolored graphs. Based on this, we prove an asymptotic formula for the number of regular edge-bicolored graphs with arbitrary weights assigned to different vertex incidence structures. The asymptotic behavior is governed by the critical points of a polynomial. As an application, we discuss the Ising model on a random 4-regular graph and show how its phase transitions arise from our formula.

1. Introduction

In this article, we study the following family of bivariate integrals,

$$I(z) = \frac{z}{2\pi} \int_{D} \exp(zg(x,y)) \, dxdy, \tag{1}$$

where D is a certain subset of \mathbb{R}^2 , g is a function $D \to \mathbb{R}$ fulfilling specific conditions (see Section 2) and z is a large positive number. Integrals as I(z) arise naturally in two important applications. First, they appear in Bayesian statistics as *marginal likelihood integrals* (see, e.g., [14, §1]). Second, they are *path integrals* associated to a zero-dimensional quantum system with two interacting

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fields parametrized by x and y, whose action is given by g(x,y) (see, e.g., [13, §2] or [2]). The setups might differ, however, in the integration domain D, leading to different asymptotic behaviors (see [12] for an asymptotic analysis in the realm of statistics). Edge-bicolored graphs play a classical role in Ramsey theory (see, e.g., [3, Ch. 12]) and their (asymptotic) enumeration is a subject with a long history (see, e.g., [15] and the references therein).

We will explain that the coefficients of the large-z asymptotic expansion of I(z) count weighted edge-bicolored graphs. Each graph is weighted by the reciprocal of the order of its automorphism group and the product of an arbitrary set of parameters assigned to each bicolored incidence structure of a vertex. We do so by proving a bivariate version of the Laplace method in Section 2, before interpreting the coefficients of the asymptotic expansion combinatorially in Section 3. Therefore, we may interpret I(z) as a generating function of edge-bicolored graphs (Theorem 3.6). From a physical perspective, these are Feynman graphs of the corresponding path integral. In Section 4, we derive an effective algorithm for the computation of those coefficients. In the final Section 5, we prove an asymptotic formula for the weighted number of regular edge-bicolored graphs, in the limit where the number of edges and vertices goes to infinity. Our main result Theorem 5.3 relates this asymptotic formula to the critical points of the polynomial g(x,y) whose shape is governed by the vertex incidence structure of the graphs. We showcase that, unlike the monochromatic case, which has previously been discussed (see, e.g., [5, Ch. 3]), only critical points satisfying some reality constraints contribute to the asymptotics.

Throughout the text we illustrate our statements through the example of the Ising model on a random 4-regular graph. The Ising model is a central object of study in mathematical physics (see, e.g., [9, 10]). The relationship between our combinatorial approach and this model is explained in Remark 3.8.

2. Laplace method and asymptotic expansions

We will start by using the Laplace method to study the large-z behavior of the integral I(z) defined in (1). We require the data D and $g: D \to \mathbb{R}$ determining I(z) to be chosen such that

- 1. the integral I(z) exists for z > 0,
- 2. D is a neighborhood of the origin,
- 3. g attains its unique supremum $\sup_{(x,y)\in D} g(x,y) = g(\mathbf{0})$ at the origin,

4. near the origin, g is analytic with absolutely converging expansion

$$g(x,y) = -\frac{x^2}{2} - \frac{y^2}{2} + \sum_{\substack{u,w \ge 0 \\ u+w \ge 3}} \Lambda_{u,w} \frac{x^u y^w}{u!w!}.$$
 (2)

The last condition ensures that (1) resembles a Gaussian integral when x and y are small. This observation allows to approximate I(z) by a slightly perturbed Gaussian when z is large.

We define a family of polynomials $a_{s,t}$ indexed by integers $s,t \geq 0$ in a twofold infinite set of variables $\lambda_{u,w}$ indexed by $u,w \geq 0$ with $u+w \geq 1$. Let \mathcal{R} be the ring of polynomials in these variables, i.e., $\mathcal{R} = \mathbb{Q}[\lambda_{0,1},\lambda_{1,0},\lambda_{1,1},\lambda_{0,2},\ldots]$. The polynomials $a_{s,t}(\lambda) \in \mathcal{R}$ are defined by the generating function

$$\sum_{s,t\geq 0} a_{s,t}(\lambda) x^s y^t = \exp\left(\sum_{\substack{u,w\geq 0\\u+w\geq 1}} \lambda_{u,w} \frac{x^u y^w}{u!w!}\right) \in \mathcal{R}[[x,y]]. \tag{3}$$

For instance, $a_{0,0}(\lambda) = 1$, $a_{1,0}(\lambda) = \lambda_{1,0}$, and $a_{2,0}(\lambda) = \frac{1}{2}(\lambda_{2,0} + \lambda_{1,0}^2)$.

We will relate the asymptotic expansion of I(z) for large z to the polynomials $a_{s,t}$. For a given function h(z), the set $\mathcal{O}(h(z))$ consists of all functions f(z) for which $\limsup_{z\to\infty}|f(z)/h(z)|$ is finite. The notation $f(z)=g(z)+\mathcal{O}(h(z))$ means that $f(z)-g(z)\in\mathcal{O}(h(z))$. The asymptotic expansion notation $f(z)\sim\sum_{n\geq 0}g_n(z)$ means that $f(z)-\sum_{n=0}^{R-1}g_n(z)\in\mathcal{O}(g_R(z))$ for all $R\geq 0$.

Proposition 2.1. If I(z), g, D and the coefficients $\Lambda_{u,w}$ are related as above, then

$$I(z) \sim \sum_{n \geq 0} A_n z^{-n},$$

for large z, where A_n is the coefficient of z^{-n} in the formal power series

$$\sum_{s,t\geq 0} z^{-(s+t)} (2s-1)!! \cdot (2t-1)!! \cdot a_{2s,2t}(z \cdot \Lambda) \in \mathbb{R}[[z^{-1}]],$$

where $(2s-1)!! = (2s-1)(2s-3)\cdots 3\cdot 1$ and $a_{2s,2t}(z\cdot \Lambda)$ is the polynomial $a_{2s,2t}(\lambda)$ defined in (3), with

$$\lambda_{u,w} = \begin{cases} 0 & \text{for } u, w \ge 0 \text{ and } 1 \le u + w < 3, \\ z\Lambda_{u,w} & \text{for } u, w \ge 0 \text{ and } u + w \ge 3. \end{cases}$$

$$(4)$$

The proof of this proposition uses the classical *Laplace method* which gives an expression for the asymptotic expansion of the integral I(z). See, e.g., [7, Appendix A] for the proof of the one-dimensional case.

Proof. Fix an integer $R \ge 0$ and any value for $\gamma \in \left(\frac{1}{3}, \frac{1}{2}\right)$. We first prove that the integral I(z) is concentrated in the square $B(z) = [-z^{-\gamma}, z^{-\gamma}]^2 \subset D$ that shrinks for growing z. Let $M(z) = \max_{(x,y) \in D \setminus B(z)} g(x,y)$, then

$$\left| I(z) - \frac{z}{2\pi} \int_{B(z)} \exp(zg(x,y)) dxdy \right| = \frac{z}{2\pi} \int_{D \setminus B(z)} \exp(zg(x,y)) dxdy$$
$$\leq \frac{z}{2\pi} \exp((z-1)M(z)) \int_{D} \exp(g(x,y)) dxdy.$$

The last integral is finite by requirement. As the origin is the unique global maximum of g in D, the maximal value M(z) will be attained on the boundary of the square B(z) if z is sufficiently large. Near the origin g(x,y) behaves as $-\frac{x^2}{2} - \frac{y^2}{2} + \text{(higher order terms)}$, so $M(z) = -\frac{1}{2}z^{-2\gamma} + \mathcal{O}(z^{-3\gamma})$. Hence,

$$I(z) = \frac{z}{2\pi} \int_{B(z)} \exp(zg(x,y)) dxdy + \mathcal{O}(z\exp(-z^{1-2\gamma})).$$
 (5)

As $\gamma < \frac{1}{2}$, we have, in particular, $\mathcal{O}(z\exp(-z^{1-2\gamma})) \subset \mathcal{O}(z^{-R})$. So, for the purpose of finding the asymptotic expansion of I(z) in decreasing powers z^0 , $z^{-1}, z^{-2}, \dots, z^{-R+1}$, integrating only over B(z) as in (5) is sufficient.

Note that, by (4), $a_{s,t}(z \cdot \Lambda)$ is a polynomial of degree at most $\frac{s+t}{3}$ in z. The function $\exp(z(\frac{1}{2}x^2 + \frac{1}{2}y^2 + g(x,y)))$ is analytic for all $(x,y) \in B(z)$. Therefore, for each K > 0, there is a constant C > 0 such that

$$\left| \exp \left(z \sum_{\substack{u,w \geq 0 \\ u+w \geq 3}} \Lambda_{u,w} \frac{x^u y^w}{u!w!} \right) - \sum_{\substack{s,t \geq 0 \\ s+t < K}} a_{s,t}(z \cdot \Lambda) x^s y^t \right| \leq C z^{\frac{1}{3}K - \gamma K} \text{ for } (x,y) \in B(z).$$

Next, we fix $K = \frac{3R}{3\gamma - 1} \ge 0$ so that $z^{\frac{1}{3}K - \gamma K} = z^{-R}$, and use (5) to get

$$I(z) = \frac{z}{2\pi} \sum_{\substack{s,t \ge 0\\ s+t < K}} a_{s,t}(z \cdot \Lambda) \int_{B(z)} e^{-z\frac{x^2}{2} - z\frac{y^2}{2}} x^s y^t dx dy + \mathcal{O}(z^{-R}).$$
 (6)

We want now to extend the integration domain to the whole real plane. For any integer $s \ge 0$, consider the integral

$$\int_{z^{-\gamma}}^{\infty} e^{-z\frac{x^2}{2}} x^s dx = \exp\left(-\frac{z^{1-2\gamma}}{2}\right) \int_0^{\infty} \exp\left(-z\frac{x^2}{2} - z^{1-\gamma}x\right) (z^{-\gamma} + x)^s dx.$$

For fixed z, the function $x \mapsto \exp(-z^{1-\gamma}x)(z^{-\gamma}+x)^s$ attains its unique maximum at $x = x_{\max} = sz^{\gamma-1} - z^{-\gamma}$. If z is sufficiently large we have $x_{\max} \le 0$.

Hence, in the range we are interested in, the integral is decreasing in x, and using $\sqrt{\frac{z}{2\pi}} \int_{\mathbb{R}} e^{-z\frac{x^2}{2}} dx = 1$ we get

$$\sqrt{\frac{z}{2\pi}} \int_{z^{-\gamma}}^{\infty} e^{-z\frac{x^2}{2}} x^s \mathrm{d}x \in \mathcal{O}\left(z^{-\gamma s} \exp\left(-\frac{z^{1-2\gamma}}{2}\right)\right) \subset \mathcal{O}(z^{-R}).$$

Combining this with (6) shows that

$$I(z) = \frac{z}{2\pi} \sum_{\substack{s,t \geq 0\\ s+t < K}} a_{s,t}(z \cdot \Lambda) \int_{\mathbb{R}^2} e^{-z\frac{x^2}{2} - z\frac{y^2}{2}} x^s y^t dx dy + \mathcal{O}(z^{-R}).$$

Using the Gaussian integral identities $\sqrt{\frac{z}{2\pi}} \int_{\mathbb{R}} e^{-z\frac{x^2}{2}} x^{2s} dx = z^{-s} \cdot (2s-1)!!$ and $\int_{\mathbb{R}} e^{-z\frac{x^2}{2}} x^{2s+1} dx = 0$ for all integers $s \ge 0$, proves the statement.

Example 2.2. Fix $D = [-1,1]^2$ and $g(x,y) = -\frac{x^2}{2} - \frac{y^2}{2} + \frac{x^4}{4!} + \lambda \frac{x^2y^2}{4} + \lambda^2 \frac{y^4}{4!}$ with $\lambda \in \mathbb{R}_{>0}$ some arbitrary positive constant. The conditions for Proposition 2.1 are fulfilled and the associated integral I(z) as defined in (1) has an asymptotic expansion $I(z) \sim \sum_{n \geq 0} A_n z^{-n}$. Using the formula from Proposition 2.1 and the generating function for the polynomials $a_{s,t}$ from (3), we find that $A_0 = 1$ and

$$A_{1} = \frac{1}{8} + \frac{1}{4}\lambda + \frac{1}{8}\lambda^{2},$$

$$A_{2} = \frac{35}{384} + \frac{5}{32}\lambda + \frac{19}{64}\lambda^{2} + \frac{5}{32}\lambda^{3} + \frac{35}{384}\lambda^{4},$$

$$A_{3} = \frac{385}{3072} + \frac{105}{512}\lambda + \frac{1295}{3072}\lambda^{2} + \frac{175}{256}\lambda^{3} + \frac{1295}{3072}\lambda^{4} + \frac{105}{512}\lambda^{5} + \frac{385}{3072}\lambda^{6}. \quad \diamond$$

In the next section, we endow the obtained analytic expressions with a combinatorial interpretation. This process is inspired by quantum field theory, where perturbative expansions of observables, which are combinatorially controlled via *Feynman graphs*, relate to *path integrals*. The integral in (1) can be seen as a path integral for a zero-dimensional space-time: the integral is then taken over all two-parameter fields on a point, hence an integral over \mathbb{R}^2 . The associated Feynman graphs are *edge-bicolored graphs*. See [13] for more details.

3. Edge-bicolored graphs

A *graph* is a one-dimensional, finite CW complex, sometimes also called multigraph in the literature. It is *edge-bicolored* if each edge has one of two different colors. We will represent graphs using only discrete data. A *(set)* partition P of a finite set H is a set of non-empty and mutually disjoint subsets of H whose union equals H. The elements of P are called *blocks*.



Figure 1: An edge-bicolored graph with two connected components.

Definition 3.1. Given two disjoint finite sets S and T of labels, an [S,T]-labeled edge-bicolored graph is a tuple $\Gamma = (V, E_S, E_T)$, where

- 1. the *vertex set V* is a partition of $S \sqcup T$,
- 2. E_S is a partition of S into blocks of size 2,
- 3. E_T is a partition of T into blocks of size 2.

We think of the elements of S and T as half-edge labels colored red and yellow, respectively. These half-edges are bundled together in vertices via the partition V. The edge sets E_S and E_T pair the half-edges into edges of the respective color. Every edge-bicolored graph without isolated vertices can be represented by at least one [S,T]-labeled graph. All graphs in this article will be edge-bicolored, so we will drop this adjective from now on.

Example 3.2. Let $S = \{s_1, s_2, \dots, s_6\}$ and $T = \{t_1, t_2\}$. The partitions

$$V = \{\{s_1, s_2, s_3, s_4\}, \{s_5, s_6, t_1, t_2\}\},\$$

$$E_S = \{\{s_1, s_2\}, \{s_3, s_4\}, \{s_5, s_6\}\},\$$

$$E_T = \{\{t_1, t_2\}\},\$$

form an [S, T]-labeled graph representing the graph depicted in Figure 1. \diamond

An *isomorphism* from an $[S_1,T_1]$ -labeled graph (V^1,E_S^1,E_T^1) to an $[S_2,T_2]$ -labeled graph (V^2,E_S^2,E_T^2) is a pair of bijections $j_S:S_1\to S_2,\ j_T:T_1\to T_2$ such that $j(V^1)=V^2,\ j(E_S^1)=E_S^2,\$ and $j(E_T^1)=E_T^2$ with j being the map that j_S and j_T canonically induce on the subsets of S,T, and $S\sqcup T$. An *automorphism* of an [S,T]-labeled graph Γ is an isomorphism to itself. Those form the group $\operatorname{Aut}(\Gamma)$.

Lemma 3.1. Each $[\{1,\ldots,2s\},\{1,\ldots,2t\}]$ -labeled graph Γ belongs to an isomorphism class of such graphs of size $\frac{(2s)!(2t)!}{|\operatorname{Aut}(\Gamma)|}$.

Proof. For given Γ , let lab(Γ) be the set of $[\{1,\ldots,2s\},\{1,\ldots,2t\}]$ -labeled graphs that are isomorphic to Γ . The group $\mathbb{S}_{2s} \times \mathbb{S}_{2t}$ acts on lab(Γ) by permuting the half-edge labels of the respective color. Aut(Γ) is the subgroup of $\mathbb{S}_{2s} \times \mathbb{S}_{2t}$ stabilizing Γ . The lemma follows from the orbit stabilizer theorem.

Example 3.3. An [S,T]-labeled graph Γ representing the graph depicted in Figure 1 has automorphism group isomorphic to $(\mathbb{S}_2 \times \mathbb{S}_2 \rtimes \mathbb{S}_2 \times \mathbb{S}_2) \times \mathbb{S}_2 \leq \mathbb{S}_6 \times \mathbb{S}_2$, where \rtimes denotes the semidirect product of groups, \mathbb{S}_6 refers to the six red halfedges in S and \mathbb{S}_2 to the two yellow half-edges in T.

We write \mathcal{G} for the set of isomorphism classes of graphs. For each $G \in \mathcal{G}$, we write $V^G, E_S^G, E_T^G, E^G = E_S^G \sqcup E_T^G$ and $\operatorname{Aut}(G)$ for the respective set or group of some [S,T]-labeled representative of G. The *Euler characteristic* of G is defined by $\chi(G) = |V^G| - |E^G|$, and does not depend on the coloring. The *bidegree* of a graph's vertex $v \in V^G$ is the pair of integers $\deg(v) = (u,w)$ where u counts the number of half-edges in v that lie in the red-colored set S and w the half-edges in the yellow-colored part T. The *vertex degree* of v is $|\deg(v)| = u + w$.

Proposition 3.4. The generating function for graphs with marked bidegrees is

$$\sum_{G\in\mathcal{G}}\frac{\eta^{|E^G|}}{|\mathrm{Aut}(G)|}\prod_{v\in V^G}\lambda_{\deg(v)}=\sum_{s,t\geq 0}\eta^{s+t}\cdot (2s-1)!!\cdot (2t-1)!!\cdot a_{2s,2t}(\lambda)\in\mathcal{R}[[\eta]],$$

where $a_{s,t}$ is defined as in (3).

We postpone the proof to after Lemma 3.2 and first illustrate the result.

Example 3.5. The formula in Proposition 3.4 provides a recipe to count our graphs for a given number of edges, grouping them according to their bidegrees. For instance, the coefficient of η^1 counts graphs with one edge:

$$\begin{split} \sum_{\substack{G \in \mathcal{G}, \\ |E^G| = 1}} \frac{1}{|\mathrm{Aut}(G)|} \prod_{\nu \in V^G} \lambda_{\deg(\nu)} = & \bigcirc + & \longleftarrow + & \bigcirc + & \longleftarrow \\ & = \frac{1}{2} \lambda_{2,0} + \frac{1}{2} \lambda_{1,0}^2 + \frac{1}{2} \lambda_{0,2} + \frac{1}{2} \lambda_{0,1}^2. \end{split}$$

Using the power series on the right-hand side of Proposition 3.4, this can be obtained simply as $a_{2,0}+a_{0,2}$, and by expanding the exponential in (3), we get exactly the above expression. If for $|E^G|=1$ these two approaches may seem equally complicated, already for graphs with two edges it is clear that the use of the generating function speeds up the computation. In fact, there are seven (monochromatic) graphs with two edges: (0,0), (0,0), (0,0), (0,0), (0,0), which turn into 23 edge-bicolored graphs. On the other hand, a simple expansion of the exponential function gives

$$\begin{split} \sum_{\substack{G \in \mathcal{G}, \\ |E^G| = 2}} \frac{1}{|\operatorname{Aut}(G)|} \prod_{\nu \in V^G} \lambda_{\deg(\nu)} &= 3a_{4,0} + a_{2,2} + 3a_{0,4} = \\ &= \lambda_{0,1} \lambda_{1,0} \lambda_{1,1} + \frac{\lambda_{0,1}^4}{8} + \frac{3\lambda_{0,1}^2 \lambda_{0,2}}{4} + \frac{\lambda_{0,1}^2 \lambda_{1,0}^2}{4} + \frac{\lambda_{0,1}^2 \lambda_{2,0}}{4} \\ &\quad + \frac{\lambda_{0,1} \lambda_{0,3}}{2} + \frac{\lambda_{0,1} \lambda_{2,1}}{2} + \frac{3\lambda_{0,2}^2}{8} + \frac{\lambda_{0,2} \lambda_{1,0}^2}{4} + \frac{\lambda_{0,2} \lambda_{2,0}}{4} + \frac{\lambda_{0,4}}{8} + \frac{\lambda_{1,0}^4}{8} \\ &\quad + \frac{3\lambda_{1,0}^2 \lambda_{2,0}}{4} + \frac{\lambda_{1,0} \lambda_{1,2}}{2} + \frac{\lambda_{1,0} \lambda_{3,0}}{2} + \frac{\lambda_{1,1}^2}{2} + \frac{3\lambda_{2,0}^2}{8} + \frac{\lambda_{2,2}}{4} + \frac{\lambda_{4,0}}{8}. \end{split}$$

To prove Proposition 3.4, we use the following lemma on the number of partitions of a set where elements come in two different colors. Let $S = \{1, ..., s\}$ and $T = \{1, ..., t\}$ and $\mathcal{P}_{s,t}$ the set of partitions of the disjoint union $S \sqcup T$. For each block B of a partition $P \in \mathcal{P}_{s,t}$ we define the *bidegree* $\deg(B)$ of the block to be the pair of integers (u, w) where u is the number of elements from S and W the number of elements from T in B.

Lemma 3.2. Given $s, t \ge 0$, consider a set of non-negative integers $n_{u,w}$ indexed by pairs u, w with $0 \le u \le s$, $0 \le w \le t$, $u + w \ge 1$, such that

$$\sum_{u} u \cdot n_{u,w} = s, \quad \sum_{w} w \cdot n_{u,w} = t.$$

The number of partitions in $\mathcal{P}_{s,t}$ with exactly $n_{u,w}$ blocks of bidegree (u,w) is

$$\frac{s!t!}{\prod_{u,w} n_{u,w}! (u!)^{n_{u,w}} (w!)^{n_{u,w}}}.$$

Proof. The group $\mathbb{S}_s \times \mathbb{S}_t$ acts on $\mathcal{P}_{s,t}$ by permuting the elements of S and T, respectively. This action is transitive if we restrict to partitions with specific block bidegree set $\{n_{u,w}\}_{u,w}$. A specific partition with given block bidegrees is stabilized by the subgroup that permutes the elements inside each block and blocks of the same size. This subgroup is isomorphic to $(\mathbb{S}_u \times \mathbb{S}_w)^{n_{u,w}} \times \mathbb{S}_{n_{u,w}}$. The claim follows from the orbit stabilizer theorem.

Proof of Proposition 3.4. From (3), and $e^X = \sum_{n \ge 0} \frac{X^n}{n!}$, we get

$$s! \cdot t! \cdot a_{s,t}(\lambda) = \sum_{\{n_{u,w}\}} \frac{s!t!}{\prod_{u,w} n_{u,w}! (u!)^{n_{u,w}} (w!)^{n_{u,w}}} \prod_{\substack{u,w \ge 0 \\ u+w > 1}} \lambda_{u,w}^{n_{u,w}}, \tag{7}$$

where the sum is over all sets of integers $\{n_{u,w}\}$ that fulfill the conditions for Lemma 3.2 with respect to s and t.

We can match the elements of the set $S = \{1, ..., 2s\}$ among each other in (2s-1)!! ways and the ones of $T = \{1, ..., 2t\}$ analogously. So, by Definition 3.1, Lemma 3.2, and (7), the number of [S,T]-labeled graphs with exactly $n_{u,w}$ vertices of bidegree u,w is $(2s-1)!! \cdot (2t-1)!!$ times the coefficient of the monomial $\prod_{u,w} \lambda_{u,w}^{n_{u,w}}$ in the polynomial $(2s)! \cdot (2t)! \cdot a_{2s,2t}(\lambda) \in \mathcal{R}$. The statement follows then from Lemma 3.1.

Our first main result follows by combining Propositions 2.1 and 3.4.

Theorem 3.6. If I(z), g, D and the coefficients $\Lambda_{u,w}$ are related as in Section 2, then the integral in (1) has the asymptotic expansion

$$I(z) \sim \sum_{n \geq 0} A_n z^{-n},$$

for large z, with the coefficients A_n given by

$$A_n = \sum_{G \in \mathcal{G}_{-n}^*} \frac{1}{|\operatorname{Aut}(G)|} \prod_{v \in V^G} \Lambda_{\deg(v)},$$

where we sum over the set \mathcal{G}_{-n}^{\star} of all isomorphism classes of edge-bicolored graphs with vertex degrees at least 3 and Euler characteristic equal to -n.

Proof. Using the fact that $\chi(G) = |V^G| - |E^G|$ we rewrite

$$\sum_{n\geq 0} A_n z^{-n} = \sum_{G\in \mathcal{G}_{-n}^{\star}} \frac{z^{\chi(G)}}{|\operatorname{Aut}(G)|} \prod_{v\in V^G} \Lambda_{\deg(v)} = \sum_{G\in \mathcal{G}_{-n}^{\star}} \frac{z^{-|E^G|}}{|\operatorname{Aut}(G)|} \prod_{v\in V^G} z \cdot \Lambda_{\deg(v)}.$$

Applying Proposition 3.4 for $\lambda_{u,w}$ as defined in (4), this is further equal to

$$\sum_{s,t>0} z^{-(s+t)} \cdot (2s-1)!! \cdot (2t-1)!! \cdot a_{2s,2t}(\lambda).$$

By Proposition 2.1, this is the large-z asymptotic expansion of I(z).

Example 3.7. Continuing Example 2.2, let $c_n^{(k)}$ be the coefficient of λ^k in A_n . By Theorem 3.6, $c_n^{(k)}$ counts automorphism-weighted graphs with Euler characteristic -n and vertex degree four, such that k_1 vertices have exactly two yellow half-edges and k_2 vertices have four yellow half-edges, so that $k_1 + k_2 = k$. We can view this explicitly for n = 2, as follows. Among the 21 (monochromatic) graphs with $\chi = -2$, there are only three 4-regular graphs. These are

$$c_2^{(0)} = \bigcirc + \bigcirc + \bigcirc + \bigcirc = \frac{1}{128} + \frac{1}{48} + \frac{1}{16} = \frac{35}{384},$$

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$$c_2^{(1)} = \bigcirc + \bigcirc = \frac{1}{32} + \frac{1}{8} = \frac{5}{32},$$

$$c_2^{(2)} = \bigcirc + \bigcirc + \bigcirc + \bigcirc = \frac{1}{64} + \frac{1}{32} + \frac{1}{8} + \frac{1}{16} + \frac{1}{16} = \frac{19}{64},$$

$$c_2^{(3)} = \bigcirc + \bigcirc = \frac{1}{32} + \frac{1}{8} = \frac{5}{32},$$

$$c_2^{(4)} = \bigcirc + \bigcirc + \bigcirc = \frac{1}{128} + \frac{1}{48} + \frac{1}{16} = \frac{35}{384}.$$

Remark 3.8 (Ising model). Our examples are motivated from the physical Ising model. The partition function of the *critical Ising model* on a specific monochromatic graph G (not necessarily lattice-like) is defined by

$$Z(G,\lambda) = \sum_{\substack{\gamma \subset G \ \gamma ext{ Eulerian}}} \lambda^{|E(\gamma)|},$$

where we sum over all Eulerian subgraphs γ of G (see, e.g., [8]). This means that if we delete all edges of G that are not in γ , then the resulting graph shall only have vertices of even degree. A pair (G, γ) of a monochromatic graph G and an Eulerian subgraph $\gamma \subset G$ is equivalent to an edge-bicolored graph in which an even number of yellow edges belongs to each vertex.

Notice that we effectively designed the polynomial g(x, y) from Example 2.2 and equivalently the coefficients $\Lambda_{u,w}$, such that the coefficient of λ^k in A_n is the automorphism-weighted number of 4-regular graphs with k yellow edges where an even number of yellow edges belong to each vertex.

Hence, with A_n as defined in Example 2.2, we find that

$$A_n = \sum_{G} \frac{Z(G, \lambda)}{|\operatorname{Aut} G|},$$

where we sum over all monochromatic graphs G that are 4-regular and which have Euler characteristic -n. We can thus interpret A_n as the partition function of the critical Ising model of a random 4-regular monochromatic graph of fixed Euler characteristic. Here, random means that each monochromatic graph G is sampled with probabilty $1/|\operatorname{Aut} G|$.

4. Efficient computation of the coefficients A_n

In this section, we describe an effective algorithm to compute the coefficients A_n that encode the asymptotic expansion of the integral (1), and the weighted numbers of edge-bicolored graphs of Euler characteristic -n, by Theorem 3.6. The algorithm is implemented in Julia and is available at [6].

Proposition 4.1. For a given integer $n \ge 1$, and the coefficients $\Lambda_{u,w}$ as required by Theorem 3.6, the following algorithm correctly computes A_0, \ldots, A_n :

Step 1: Define the polynomials

$$F_k(x,y) = \sum_{\substack{u,w \ge 0 \\ u+w=k+2}} \Lambda_{u,w} \frac{x^u y^w}{u!w!} \text{ for } k \in \{1,\dots,2n\};$$

Step 2: Set $Q_0(x,y) = 1$ and recursively compute Q_1, \dots, Q_{2n} using

$$Q_m(x,y) = \frac{1}{m} \sum_{k=1}^m k F_k(x,y) Q_{m-k}(x,y) \text{ for } m \in \{1,\dots,2n\};$$

Step 3: Let $q_{s,t}^{(k)}$ be the coefficients of $Q_k(x,y) = \sum_{s,t \ge 0} q_{s,t}^{(k)} x^s y^t$. Then,

$$A_k = \sum_{s,t \ge 0} (2s-1)!! \cdot (2t-1)!! \cdot q_{2s,2t}^{(2k)} \text{ for } k \in \{0,\dots,n\}.$$

To run the algorithm with a fixed n, it is sufficient to know $\Lambda_{u,w}$ for all $u, w \ge 0$ with $u + w \le 2n + 2$. Also, recall that we require $\Lambda_{u,w} = 0$ if u + w < 3.

Proof. By Proposition 2.1, we have this identity of power series in z^{-n} :

$$\sum_{n>0} A_n z^{-n} = \sum_{s,t>0} z^{-(s+t)} \cdot (2s-1)!! \cdot (2t-1)!! \cdot a_{2s,2t}(z \cdot \Lambda),$$

where $a_{s,t}(z \cdot \Lambda)$ is as described in Proposition 2.1 and (3):

$$\sum_{s,t\geq 0} a_{s,t}(z\cdot \Lambda)x^s y^t = \exp\left(z\sum_{\substack{u,w\geq 0\\u+w\geq 3}} \Lambda_{u,w} \frac{x^u y^w}{u!w!}\right).$$

Rescaling $(x,y) \mapsto (x/\sqrt{z},y/\sqrt{z})$ in the above formula gives the following identity of power series in $\mathbb{R}[x,y][[z^{-1/2}]]$,

$$\sum_{s,t\geq 0} z^{-\frac{s+t}{2}} a_{s,t}(z \cdot \Lambda) x^s y^t = \exp\left(\sum_{k\geq 1} z^{-\frac{k}{2}} F_k(x,y)\right),\tag{8}$$

where we used the definition of $F_k(x,y)$ in the statement. Let $q_{s,t}^{(k)}$ be the coefficients $\sum_{k\geq 0}q_{s,t}^{(k)}z^{-\frac{k}{2}}=z^{-\frac{s+t}{2}}a_{s,t}(z\cdot\Lambda)$. With this definition, (8) and the formula under Step 3 in the statement correctly compute A_k .

It remains to prove that the coefficients $q_{s,t}^{(k)}$ are computed correctly by Step 2 in the statement. Rewrite (8) using the definition of $Q_k(x,y)$, before applying the derivative operator $z\frac{\partial}{\partial z}$ on both sides. This gives

$$\begin{split} z\frac{\partial}{\partial z}\left(\sum_{k\geq 0}Q_k(x,y)z^{-\frac{k}{2}}\right) &= z\frac{\partial}{\partial z}\exp\left(\sum_{k\geq 1}z^{-\frac{k}{2}}F_k(x,y)\right)\\ \Rightarrow &-\sum_{k>0}\frac{k}{2}Q_k(x,y)z^{-\frac{k}{2}} = -\left(\sum_{m>0}Q_m(x,y)z^{-\frac{m}{2}}\right)\sum_{k>1}\frac{k}{2}z^{-\frac{k}{2}}F_k(x,y). \end{split}$$

The recursive relation between Q_m and F_k follows by comparing the $z^{-\frac{m}{2}}$ coefficients on both sides of this equation.

5. Asymptotics and critical points

In this section, we study the asymptotic behavior of the coefficients A_n in Theorem 3.6 for large n. Here, we will restrict ourselves to *regular* edge-bicolored graphs, meaning that each vertex has a fixed degree $k \ge 3$. For fixed coefficients $\Lambda_{u,w}$ given for $u,w \ge 0$ with u+w=k, we study the weighted sum over graphs

$$A_n = \sum_{G \in \mathcal{G}_{-n}^k} \frac{1}{|\mathrm{Aut}(G)|} \prod_{v \in V^G} \Lambda_{\deg(v)},$$

where \mathcal{G}_{-n}^k is the set of all regular (edge-bicolored) graphs with vertex degree k and Euler characteristic $-n = |V^G| - |E^G|$. As for each k-regular graph G we have $k|V^G| = 2|E^G|$, all graphs in \mathcal{G}_{-n}^k have $\frac{2n}{k-2}$ vertices and $\frac{nk}{k-2}$ edges. It is convenient to define the homogeneous polynomial

$$V(x,y) = g(x,y) + \frac{x^2}{2} + \frac{y^2}{2} = \sum_{\substack{u,w \ge 0 \\ u+w=k}} \Lambda_{u,w} \frac{x^u y^w}{u!w!} \in \mathbb{R}[x,y].$$

Let Φ be the set of global maxima of the function

$$S^1 = \{(x, y) \in \mathbb{R}^2 : x^2 + y^2 = 1\} \to \mathbb{R}_{\geq 0}, \quad (x, y) \mapsto |V(x, y)|.$$

A point $(x,y) \in \Phi$ is non-degenerate if $k^2V(x,y) \neq \left(\frac{\partial^2 V}{\partial x^2}(x,y) + \frac{\partial^2 V}{\partial y^2}(x,y)\right)$.

Proposition 5.1. Let $M = \frac{k}{k-2}$ and $K = \frac{2}{k-2}$. If $A_n, V, \Lambda_{u,w}$ and Φ are related as described above and all extrema in Φ are non-degenerate, then

$$A_n \sim \begin{cases} \frac{1}{2\sqrt{2}\pi} k^{nM+\frac{1}{2}} K^{n-\frac{1}{2}} \Gamma(n) \sum\limits_{(x,y) \in \Phi} \frac{V(x,y)^{nK}}{\sqrt{B(x,y)}} & \text{if } nK, nM \in \mathbb{Z}, \\ 0 & \text{else,} \end{cases}$$

where Γ denotes the Gamma function and

$$B(x,y) = k^2 - \frac{\frac{\partial^2 V}{\partial x^2}(x,y) + \frac{\partial^2 V}{\partial y^2}(x,y)}{V(x,y)} \quad \text{for } (x,y) \in S^1.$$

We will prove this theorem by first proving an integral representation of the coefficients A_n . Afterwards, we apply the one-dimensional Laplace method to provide an asymptotic expression for this integral in the large n limit.

Lemma 5.1. Let $M = \frac{k}{k-2}$ and $K = \frac{2}{k-2}$. For a given integer $n \ge 0$ such that nK and nM are integers, we have

$$A_n = \frac{2^{nM}(nM)!}{2\pi \cdot (nK)!} \int_{-\pi}^{\pi} V(\cos\varphi, \sin\varphi)^{nK} d\varphi.$$

If nK or nM is not an integer then $A_n = 0$.

Proof. A k-regular graph has nM edges and nK vertices, so nM and nK must be integers; otherwise $A_n = 0$. We will assume the former. By Proposition 3.4,

$$A_n = \sum_{\substack{s,t \ge 0\\ s+t = nM}} (2s-1)!! \cdot (2t-1)!! \cdot a_{2s,2t}(\Lambda).$$

If s+t=nM, it follows from (3) that $a_{2s,2t}(\Lambda)$ is a homogeneous polynomial of degree nK. Because $\exp(X)=\sum_{N\geq 0}\frac{X^N}{N!}$, it also follows that $a_{2s,2t}(\Lambda)$ is the coefficient in front of $x^{2s}y^{2t}$ in the quotient $V(x,y)^{nK}/(nK)!$, since V is homogeneous. Using $\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}}e^{-\frac{x^2}{2}}x^{2s}\mathrm{d}x=(2s-1)!!$ and $\int_{\mathbb{R}}e^{-\frac{x^2}{2}}x^{2s+1}\mathrm{d}x=0$ for integers s, we obtain

$$A_n = \frac{1}{2\pi \cdot (nK)!} \int_{\mathbb{R}^2} e^{-\frac{x^2}{2} - \frac{y^2}{2}} V(x, y)^{nK} dx dy.$$

We can pass to polar coordinates and use $V(rx, ry) = r^k V(x, y)$ together with

$$\int_0^\infty e^{-\frac{r^2}{2}} r^{nkK+1} dr = \int_0^\infty e^{-q} (2q)^{nkK/2} dq = 2^{nM} (nM)!$$

to prove the lemma.

Proof of Proposition 5.1. We are interested in the cases in which $A_n \neq 0$. When n is large, the main contribution to the integral in the statement of Lemma 5.1 comes from angles φ where $|V(\cos\varphi,\sin\varphi)|$ is maximal. Let φ_c be the location of such a maximum. By definition, we have $(\cos\varphi_c,\sin\varphi_c) \in \Phi$. Near this maximum, we get the Taylor expansion

$$f_{oldsymbol{arphi}_c}(oldsymbol{arphi}) := \log rac{V(\cos oldsymbol{arphi}, \sin oldsymbol{arphi})}{V(\cos oldsymbol{arphi}_c, \sin oldsymbol{arphi}_c)} = -B(\cos oldsymbol{arphi}_c, \sin oldsymbol{arphi}_c) rac{(oldsymbol{arphi} - oldsymbol{arphi}_c)^2}{2} + \mathcal{O}((oldsymbol{arphi} - oldsymbol{arphi}_c)^3),$$

where $B(\cos \varphi_c, \sin \varphi_c)$ is defined as in the statement. Because φ_c is a maximum of $|V(\cos \varphi_c, \sin \varphi_c)|$, we have $B(\cos \varphi_c, \sin \varphi_c) \ge 0$. Our assumption that all the maxima are non-degenerate hence implies that $B(\cos \varphi_c, \sin \varphi_c) > 0$.

We may therefore write, for some sufficiently small $\varepsilon > 0$,

$$A_n = \frac{2^{nM}(nM)!}{2\pi \cdot (nK)!} \left(\sum_{(\cos \varphi_c, \sin \varphi_c) \in \Phi} V(\cos \varphi_c, \sin \varphi_c)^{nK} \int_{\varphi_c - \varepsilon}^{\varphi_c + \varepsilon} e^{nKf_{\varphi_c}(\varphi)} d\varphi \right) + R_1(n, \varepsilon).$$

From the Taylor expansion of the function $f_{\varphi_c}(\varphi)$ and Lemma 5.1, it follows by the same reasoning as in the proof of Proposition 2.1 that the remainder term respects the bound $|R_1(n,\varepsilon)| \leq C_1 \exp(-C_2 n \varepsilon^2)$ with some constants $C_1, C_2 > 0$. Specifying $\varepsilon = n^{-\gamma}$ with $\gamma \in (\frac{1}{3}, \frac{1}{2})$ allows us to truncate the Taylor expansion

of f_{φ_c} after the second term without changing asymptotic behavior in the $n \to \infty$ limit. Hence,

$$\int_{\varphi_c-\varepsilon}^{\varphi_c+\varepsilon} e^{nKf_{\varphi_c}(\varphi)} d\varphi = \int_{-\varepsilon}^{\varepsilon} \exp\left(-nK\frac{B(\cos\varphi_c,\sin\varphi_c)}{2}\varphi^2\right) d\varphi + R_2(n,\varepsilon).$$

The remainder term fulfills $|R_2(n,\varepsilon)| < C_3 n^{\frac{1}{2}} \varepsilon^3$ for some $C_3 > 0$. Again, as in the proof of Proposition 2.1, we may complete the Gaussian integral to find that

$$\int_{\varphi_c-\varepsilon}^{\varphi_c+\varepsilon} e^{nKf_{\varphi_c}(\varphi)} d\varphi = \sqrt{\frac{2\pi}{nK \cdot B(\cos\varphi_c,\sin\varphi_c)}} + \mathcal{O}(n^{-1}).$$

The result follows from Stirling's formula $\Gamma(n) \sim \sqrt{2\pi n^{-1}} n^n e^{-n}$ as $n \to \infty$.

Example 5.2. To continue the running example of the Ising model (cf. Example 3.7), let $V(x,y) = \frac{x^4}{4!} + \lambda \frac{x^2 y^2}{4} + \lambda^2 \frac{y^4}{4!}$. We want to find the critical points of V on the circle, that means the points $(x,y) \in \mathbb{R}^2$ with $x^2 + y^2 = 1$ satisfying

$$y\frac{\partial V}{\partial x}(x,y) = x\frac{\partial V}{\partial y}(x,y). \tag{9}$$

We get the following eight critical points:

$$(\pm 1, 0), (0, \pm 1), \left(\pm \frac{\sqrt{\lambda(\lambda - 3)}}{\sqrt{\lambda^2 - 6\lambda + 1}}, \pm \frac{\sqrt{1 - 3\lambda}}{\sqrt{\lambda^2 - 6\lambda + 1}}\right) \in \mathbb{R}^2. \tag{10}$$

Our case of interest is $\lambda > 0$. Then, the last four points are real if and only if $\lambda \in [\frac{1}{3},3]$, and in that interval those are the maxima of |V| on S^1 . For $\lambda < \frac{1}{3}$, the maxima are $(\pm 1,0)$, whereas for $\lambda > 3$, the maxima are $(0,\pm 1)$. Figure 2 displays the function (V(x,y)x,V(x,y)y) and its critical points, for $\lambda \in (0,4)$.

We can now use Proposition 5.1 to find $A_n \sim c \Gamma(n)\alpha^n$, where $c = c(\lambda)$ and $\alpha = \alpha(\lambda)$ are piecewise defined as

$$\begin{array}{c|cccc} & \alpha(\lambda) & c(\lambda) \\ \hline 0 < \lambda < \frac{1}{3} & \frac{2}{3} & \frac{1}{\pi} \sqrt{\frac{1}{2 - 6\lambda}} \\ \frac{1}{3} < \lambda < 3 & \frac{-16\lambda^2}{3\lambda^2 - 18\lambda + 3} & \frac{1}{\pi} \sqrt{\frac{8\lambda}{-3\lambda^2 + 10\lambda - 3}} \\ \lambda > 3 & \frac{2\lambda^2}{3} & \frac{1}{\pi} \sqrt{\frac{\lambda}{2\lambda - 6}} \end{array}$$

The function α is continuous, it is not C^1 -differentiable at $\lambda = \frac{1}{3}$, and it is C^1 -but not C^2 -differentiable at $\lambda = 3$. On the other hand, the limits of $c(\lambda)$ at $\frac{1}{3}$,

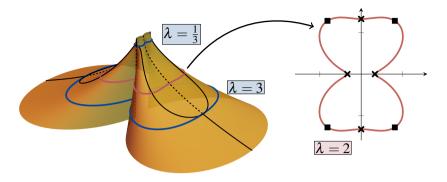


Figure 2: The system (9) for the function V from Example 5.2. Left: all values of $\lambda \in (0,4)$ on the (reversed) vertical axis. At each level $\lambda = \text{const.}$ the black continuous curves are the maxima; the dashed curves are the minima. In blue, the curves for $\lambda = \frac{1}{3}$ and $\lambda = 3$ where the behavior of the maxima changes. Right: the section $\lambda = 2$, with its maxima (squares) and its minima (crosses).

3 go to infinity from both sides. This can be observed in Figure 3. The points $\lambda = \frac{1}{3}$ and $\lambda = 3$ where the functions $\alpha(\lambda)$ and $c(\lambda)$ are non-analytic are *phase transition* points. Phase transitions are of pivotal interest in statistical physics. Here, we find the phase transitions of the Ising model on a random 4-regular graph. In each of the three regions for the parameter λ , the statistical system is expected to exhibit intrinsically different behaviors.

Note that using Proposition 4.1 we can also compute A_n for large n and solve for α and c numerically. For details see our implementation at [6].

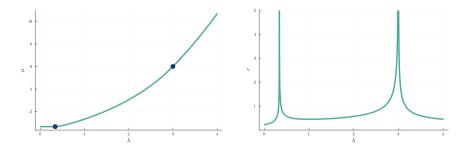


Figure 3: The behavior of $\alpha(\lambda)$ and $c(\lambda)$ in the Ising model from Example 5.2. The phase transitions at $\lambda = \frac{1}{3}$, 3 can be detected in both quantities. At $\lambda = 3$, the function α is C^1 - but not C^2 -differentiable.

It is common belief in physics that the asymptotic behavior of A_n depends on the critical points of $g(x,y) = -\frac{x^2}{2} - \frac{y^2}{2} + V(x,y)$ (see, e.g., [11]). Moreover,

it is well-known, also in applied mathematics, that identifying the critical point which contributes most to the asymptotics is a complicated *connection problem* [1]. With this in mind, we rephrase Proposition 5.1 in terms of the critical points of g instead of those of V restricted to the sphere. We write $\operatorname{crit}_D f$ for the set of critical points of f restricted to the domain f. Let

$$\Psi = \left\{ (w, z) \in \operatorname{crit}_{\mathbb{C} \cdot \mathbb{R}^2} g \setminus \{\mathbf{0}\} : \|(w, z)\| \le \|(w', z')\| \ \forall (w', z') \in \operatorname{crit}_{\mathbb{C} \cdot \mathbb{R}^2} g \setminus \{\mathbf{0}\} \right\},$$

where $\mathbb{C} \cdot \mathbb{R}^2$ is the set of complex points (w,z) whose ratio (when well-defined) is real. We call points in Ψ *non-degenerate* if the Hessian (the matrix of second derivatives) of g has full rank.

Theorem 5.3. Assume that A_n , g, and Ψ are related as described above and all extrema in Ψ are non-degenerate. Then

$$A_n \sim \frac{1}{2\pi} \Gamma(n) \sum_{(w,z) \in \Psi} \frac{(-g(w,z))^{-n}}{\sqrt{-\det H_g(w,z)}}.$$
 (11)

Proof. Our goal is to express A_n from Proposition 5.1 in terms of the critical points of g. The first step is to associate the critical points of V to those of g. Given $(x,y) \in \operatorname{crit}_{S^1}(V)$, we look for some $\ell \in \mathbb{C}^*$ such that $(\ell x, \ell y) \in \operatorname{crit}_{\mathbb{C} \cdot \mathbb{R}^2}(g)$. Imposing the conditions $\ell x = \frac{\partial V}{\partial x}(\ell x, \ell y)$, $\ell y = \frac{\partial V}{\partial y}(\ell x, \ell y)$, and using homogeneity of V, we get

$$\ell^{2-k} = kV(x, y). \tag{12}$$

Therefore, as $k \ge 3$,

$$\max_{(x,y)\in\operatorname{crit}_{S^1}(V)} V(x,y) = \frac{1}{k} \left(\min_{(w,z)\in\operatorname{crit}_{\mathbb{C}.\mathbb{R}^2}(g)} \|(w,z)\| \right)^{2-k},$$

so every element $(w, z) \in \Psi \subset \mathbb{C} \cdot \mathbb{R}^2$ arises as $(\ell x, \ell y)$ for some $(x, y) \in \Phi$.

Using these considerations, we write the result from Proposition 5.1 in terms of the critical points of g. At a point $(w,z) = (\ell x, \ell y) \in \Psi$, by (12), we have

$$g(w,z) = -\frac{\ell^2}{2} + V(w,z) = -\frac{\ell^2}{2} + \ell^k V(x,y) = \ell^2 \frac{2-k}{2k}.$$
 (13)

Let $K = \frac{2}{k-2}$ and $M = \frac{k}{k-2}$. Then, we have

$$V(x,y)^{nK} = k^{-nK}\ell^{-2n} = k^{-nK}(-kKg(w,z))^{-n} = k^{-nM}K^{-n}(-g(w,z))^{-n}.$$

This allows to cancel prefactors in the asymptotic expression for A_n from Proposition 5.1. We are left to rewrite B in terms of (w, z). Notice that the determinant

of the Hessian of g(w,z) can be expressed, using (12), as

$$\begin{split} \det & H_g(w,z) = \frac{1}{\ell^2} \det & H_g(\ell x, \ell y) = B(x,y) \ell^{k-2} V(x,y) (1 - k(k-1) \ell^{k-2} V(x,y)) \\ & = - B(x,y) \frac{k-2}{k}, \end{split}$$

where (ℓ, x, y) are new coordinates on $\mathbb{C}^* \times S^1$, and $(w, z) \in \Psi$. Hence,

$$\begin{split} A_n &\sim \frac{1}{2\sqrt{2}\pi} k^{nM + \frac{1}{2}} K^{n - \frac{1}{2}} \Gamma(n) \sum_{(x,y) \in \Phi} \frac{V(x,y)^{nK}}{\sqrt{B(x,y)}} \\ &= \frac{1}{4\pi} \sqrt{k(k-2)} \Gamma(n) \frac{2}{k-2} \sum_{(w,z) \in \Psi} \frac{(-g(w,z))^{-n}}{\sqrt{-\frac{k \det H_g(w,z)}{k-2}}} \\ &= \frac{1}{2\pi} \Gamma(n) \sum_{(w,z) \in \Psi} \frac{(-g(w,z))^{-n}}{\sqrt{-\det H_g(w,z)}}, \end{split}$$

where the factor $\frac{2}{k-2}$ appears since each of the k-2 points $\{(w,z)=(\ell x,\ell y)\}$ in Ψ is counted twice by the corresponding points $\{(x,y),(-x,-y)\}\in\Phi$.

Remark 5.4. The formula (11) yields 0 if nM or nK are not integers. Indeed, using (13) from the proof above, we can write, for $(w, z) \in \Psi$,

$$g(w,z) = (l \cdot \zeta_i)^2 \frac{2-k}{2k}, \quad i \in \{1,...,k-2\},$$

where $l \in \mathbb{R}$ and ζ_i is a (k-2)th root of unity, so $(w,z) = (l\zeta_i x, l\zeta_j y)$ for some $(x,y) \in \Phi$. Also $(l\zeta_j x, l\zeta_j y) \in \Psi$ for all $j \in \{1, \dots, k-2\}$. Therefore, the sum in (11) becomes

$$\begin{split} \sum_{(w,z)\in\Psi} \frac{(-g(w,z))^{-n}}{\sqrt{-\det H_g(w,z)}} &\propto \sum_{j=1}^{k-2} \frac{(l\cdot\zeta_j)^{-2n}\left(\frac{k-2}{k}\right)^{-n}}{\sqrt{-\det H_g(l\zeta_jx,l\zeta_jy)}} \\ &= \begin{cases} (k-2)\frac{l^{-2n}\left(\frac{k-2}{k}\right)^{-n}}{\sqrt{-\det H_g(lx,ly)}} & \text{if } (k-2) \mid 2n, \\ 0 & \text{else.} \end{cases} \end{split}$$

The condition $(k-2) \mid 2n$ is equivalent to $nK \in \mathbb{Z}$, which also implies $nM \in \mathbb{Z}$.

We exhibit the connection between the two collections Φ and Ψ of critical points explicitly in our running example.

Example 5.5. Let $V(x,y) = \frac{x^4}{4!} + \lambda \frac{x^2 y^2}{4} + \lambda^2 \frac{y^4}{4!}$ and $g(x,y) = -\frac{x^2}{2} - \frac{y^2}{2} + V(x,y)$. Consider the system of critical equations for g

$$w = \frac{\partial V}{\partial w}(w, z), \quad z = \frac{\partial V}{\partial z}(w, z), \tag{14}$$

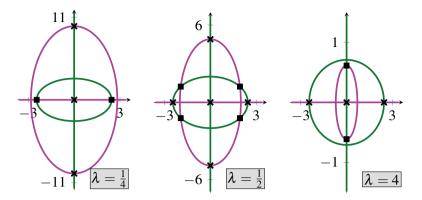


Figure 4: The system (14) for the function V from Example 5.5, for the values $\lambda = \frac{1}{2}, \frac{1}{4}, 4$, from left to right. The solutions are marked in black. The solutions that are (equally) closest to (but distinct from) the origin, are marked with squares. Notice the different scaling in the y-axis, for the sake of clarity.

and its complex non-trivial solutions, for $\lambda > 0$:

$$(\pm\sqrt{6},0),\left(0,\pm\frac{\sqrt{6}}{\lambda}\right),\left(\pm\sqrt{\frac{9-3\lambda}{4\lambda}},\pm\frac{\sqrt{9\lambda-3}}{2\lambda}\right).$$

Among these solutions, some are real for every $\lambda > 0$. The last type of singular points is real if and only if $\lambda \in [\frac{1}{3}, 3]$. We get

$$\Psi = \begin{cases} (\pm\sqrt{6},0) & 0 < \lambda < \frac{1}{3}, \\ \left(\pm\sqrt{\frac{9-3\lambda}{4\lambda}}, \pm\frac{\sqrt{9\lambda-3}}{2\lambda}\right) & \frac{1}{3} < \lambda < 3, \\ \left(0,\pm\frac{\sqrt{6}}{\lambda}\right) & \lambda > 3. \end{cases}$$

This is displayed in Figure 4. The reader may check that rescaling each point in Ψ to unit vector gives precisely two of the points in (10).

Remark 5.6. Lee–Yang theory studies the location of the roots of the polynomials A_n , when n becomes large. This fascinating theory touches combinatorics, statistics and physics (see, e.g., [4] for an overview). In the spirit of Lee–Yang theory, the two phase transitions $\lambda = \frac{1}{3}$, 3 in the running example can be detected also by looking at the asymptotic behavior of the roots of $A_n(\lambda)$ as $n \to \infty$. Using our algorithm from Proposition 4.1, we can compute the polynomials $A_n(\lambda)$ and find their roots numerically. This is the content of Figure 5. The roots of these polynomials are all complex (except for $\lambda = -1$, for odd n) but they get closer and closer to the real values $\lambda = \frac{1}{3}$ and $\lambda = 3$.

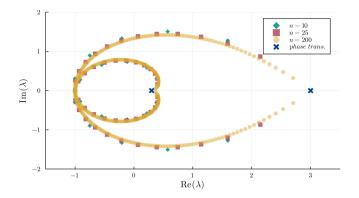


Figure 5: The roots of $A_n(\lambda)$ under the assumptions of Example 2.2, for n = 10,25,200. The blue crosses are the phase transitions $\lambda = \frac{1}{3},3$.

Although Proposition 5.1 and Theorem 5.3 assume *V* to be homogeneous, the following example shows that this condition does not seem to be necessary.

Example 5.7. Take the inhomogeneous polynomial $V(x,y) = \frac{x^3}{3!} + \lambda \frac{xy^2}{2} + \lambda^2 \frac{y^4}{4!}$, with $\lambda > 0$. For $\lambda < \frac{1}{2}$, one can compute that $\Psi = \{(2,0)\}$; for $\lambda > \frac{1}{2}$, one gets

$$\Psi = \left\{ \left(\frac{4\lambda - \sqrt{2\lambda(8\lambda - 3)}}{\lambda}, \pm \sqrt{6}\sqrt{\frac{1 - (4\lambda - \sqrt{2\lambda(8\lambda - 3)})}{\lambda^2}} \right) \right\}.$$

The formula for A_n from Theorem 5.3 would give

	$\alpha(\lambda)$	$c(\lambda)$
$0 < \lambda < \frac{1}{2}$	$\frac{3}{2}$	$\frac{1}{2\pi\sqrt{1-2\lambda}}$
$\lambda>rac{1}{2}$	$\frac{6\lambda^2}{(8\lambda - 3)\left(16\lambda - 3 - 4\sqrt{2\lambda(8\lambda - 3)}\right)}$	$\frac{1}{\pi}\sqrt{\frac{\lambda}{32\lambda^2 - 12\lambda + 2\sqrt{2\lambda(8\lambda - 3)(1 - 16\lambda^2)}}}$

This matches our numerical computations, see [6].

Based on the previous example and similar computations, we conjecture that Theorem 5.3 is also valid for inhomogeneous V(x,y), i.e., graphs that are not necessarily regular. In this setting, the univariate Laplace method as used in the proof of Proposition 5.1 does not work anymore, also due to the failure of Lemma 5.1. Instead a multivariate saddle point method shall be required.

Conjecture 5.8. Let g and A_n be related as in the beginning of Section 2 (i.e., $g(x,y) + \frac{x^2}{2} + \frac{y^2}{2}$ is not necessarily homogeneous). Let Ψ be defined as before

and assume that all points in Ψ are non-degenerate. Then,

$$A_n \sim \frac{1}{2\pi} \Gamma(n) \sum_{(w,z) \in \Psi} \frac{(-g(w,z))^{-n}}{\sqrt{-\det H_g(w,z)}}$$
 as $n \to \infty$.

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M. BORINSKY

ETH Institute for Theoretical Studies & Perimeter Institute e-mail: michael.borinsky@eth-its.ethz.ch

C. MERONI

ETH Institute for Theoretical Studies e-mail: chiara.meroni@eth-its.ethz.ch

M. WIESMANN

Max Planck Institute for Mathematics in the Sciences e-mail: wiesmann@mis.mpg.de